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BROADSIDE SCATTERING OF A TUBULAR CYLINDER FOR EVALUATION OF TARGET IDENTIFICATION

by

Boaz Haklay

March 1985

Thesis Advisor:

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Broadside Scattering of a Tubular Cylinder for Evaluation of Target Identification

bу

Boaz Haklay Lieutenant,Israeli Navy B.S., Jerusalem College of Technology, 1980

Submitted in partial fulfillment of the requirements for the degree of

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ABSTRACT

The concept of target identification through its back scattering cross section, in the resonance region was investigated. Measurements from the broadside aspect angle of several scaled tubular cylinders have been used for this purpose. The experimental results and theoretical approximation for some of the cylinders are presented. That data will serve as the baseline for further investigation in this project.

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I. INTRODUCTION

Target identification is desirable if different actions are to be taken toward different targets. Today when missiles can be sent long before a target comes within visual range, there is a need of identify target by some means other then visual.

The use of radar provides a greater range for target detection. Radars have the ability to detect the presence of a target, and to obtain information about target position, speed and acceleration. A trained operator can identify the kind of target infront of him by the size of the spot on the screen and from the direction such a target is approaching, but he may not always be correct.

Detection by a radar is done by detecting the back scattered electromagnetic signal returned from a target. The back scattered signal originates from the surface currents excited on the object when it is irradiated by an incident electromagnetic wave. Different targets have different outer surface that cause different surface currents to flow on the target when identical incident waves are irradiating the targets.

Assume the presence of a continuous incident electromagnetic wave. The incident wave keeps impinging on the target and excites new surface current which add vectorally to the existing ones. The total surface current, and therefore the back scattered signal are functions of the target shape and the wavelength of the incident wave.

Assuming a uniform plane incident wave, the back scattering cross section of a target is a quantitative measure of the ratio of power density in the vector signal scattered in the direction of the receiver to the power density of the

electromagnetic wave incident upon the target [Ref. 1], The back scattering cross section of a target as a function of the wavelength or the frequency of the incident wave can be used as a working tool for target identification

One approach to this problem, is by looking into the Rayleigh region [Ref. 2], where the frequency rangs from zero up to a wavelength about half the linear dimension of the target. A different approach is to study the scattering of a target in the resonance region where its cross section varies rapidly with frequency, and is a critical function of the shape of the target [Ref. 3]. Even though some resonances may not be observable at some particular target aspect angles. The advantages of using the resonance region are that the scattered fields of interests are stronger and thus are easier to be detected. Because the resonance frequencies depend critically on target shapes, and because there are only a finite number of targets of interest only a few resonances will be needed before a target can be identified.

This thesis is part of an ongoing project at the Naval Postgraduate School on target identification. It studies the broadside scattering of a tubular cylinder of finite length, made of very thin brass walls. The targets used in this research were a set of tubular circular cylinders with different lengthes and diameters. This shape has been chosen because of its resemblance to a missile body because its theoretical solution is available. should serve as the basis for further efforts in developing a target identification scheme employing the frequency dependence of the back scattered field from a target in the resonance region. The effects of adding fins to the cylinders will be investigated next and will lead to the last step where models of real targets will be identify through this scheme.

Chapter II describes the general approach to the solution of the back scattering cross section problem. The exact solution for the scattering by a perfectly conducting tubular cylinder having very thin walls from the broadside aspect angle is presented. Chapter III contains the experimental setup at the Naval Postgraduate School, the measurement procedure and the experimental data. Chapter presents data analysis of the results as well as comparison between theoretical data obtained from the solution shown in Chapter II. Because this was the first time frequency dependence of the cross section of tubular cylinder was ever tested at the resonance region, the results are very surprising and can be used to direct further efforts in this project. Chapter V contains a summary of the result, the problems encountered and areas of future work.

II. ANALYTICAL SOLUTION FOR SCATTERING PROBLEM

A. GENERAL APPROACH TO SCATTERING PROBLEM

A parameter that defines the scattering efficiency of a target was sought in the earliest days of radar. This parameter when refers to the equivalent isotropic reflector is called cross section and denoted by the symbol σ . The theoretical definition of back scattering cross section is given by 2.1 .

$$\sigma = 4\pi r^2 \lim_{r \to \infty} |E_s/E_i|^2 \tag{2.1}$$

Where E_i-Magnitude of electric-field component of incident electromagnetic (EM) field at the target.

 ${\sf E_s}{\sf -Magnitude}$ of electric-field component of scattered EM field as measured by a hypothetical observer

r-Distance from target to the hypothetical observer

In this definition the incident and scattered fields are introduced. The incident field means the field maintained by the oscillating charges and currents in the driving or primary antenna constitutes the source. When the target is irradiated with an incident EM wave, surface current is excited on the target. This surface current will radiate and generate the scattered field. The configuration of a scattering problem is shown in Figure 2.1.

The limiting process is introduced to assure that the distance at which the hypothetical observation is made is

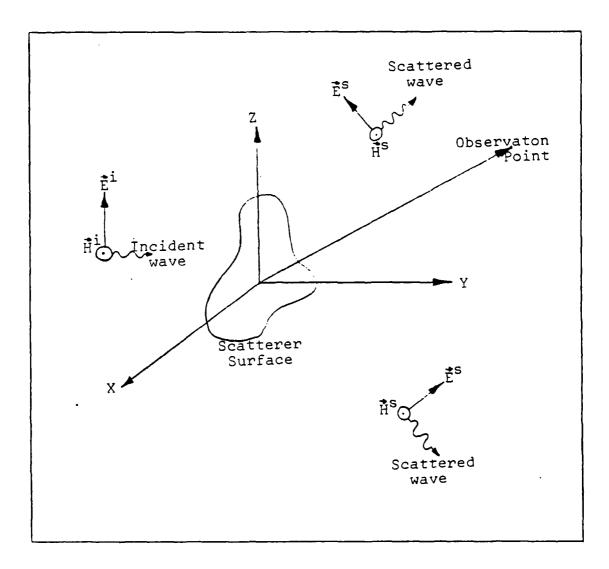


Figure 2.1 Configuration of a Scattering Problem

far enough from the target so that only the R⁻¹ dependent term of the scattered field is retained. Under the free-space conditions assumed here, the ratio $|E_{\rm s}/E_{\rm i}|^2$ is the same as the ratio of the power flux density of the scattered waves at the observer to that of the incident wave at the target.

The incident field is not independent of the presence of the target because of the coupling between the currents and charges in the source and the currents and charges on the target. To simplify the problem, it is usually assumed that the source is separated from the target by a large distance and thus the incident field is independent of the presence of the target.

Ideally, one would compute the radar cross section of a target through the formal solution of Maxwell equations under the boundary conditions appropriate to the body. The integral equation formulation shows that electromagnetic scattering of an incident wave by an arbitrary body can be described in terms of an integral of various vector products involving the surface electric and magnetic fields on the target. The Chu-Stratton integral [Ref. 4 p. 464] is convenient for this purpose. This integral is an exact representation of the scattered electromagnetic field in terms of an integration over a complete surface enclosing the body in question. In particular, if there is available knowledge of the total distribution of electric and magnetic fields about the body (or what is equivalent, surface currents), tion of these values in the Chu-Stratton integral would permit the immediate evaluation of the scattered fields. Approximate numerical solutions of such integral equations require the use of high speed digital computers to estimate surface currents flowing on the body. Because of the limitation of computation time and storage capacity, this method is applicable only when the dimensions of the target do not exceed a very few wavelengths.

Exact solutions for the scattering problem are rare. For many practical problems, only approximate solutions are obtainable. Aside from different numerical schemes, asymptotic techniques are also used when the target dimensions are much larger than the wavelength. It is called geometrical diffraction theory and it combines the simplicity inherent in the ray optics with the necessary consideration

of wavelengths and phases. This method uses the concept of scattering centers and localizes them at the points of surface discontinuity in the belief that specular diffractions occur only at these discontinuities with contributions from surface regions a few wavelength within these points. Each center is assigned a magnitude and a phase based upon asymptotic expansion of the exact solution of a two dimensional case of a plane geometry. By concentrating on scattering centers, it is possible to predict the polarization of the signal reradiated from each center and so to preserve polarization dependence in computed result. It is also possible to predict the dependence of radar scattering upon bistatic angle, but this approach cannot include resonances of the target.

B. SCATTERING BY A TUBULAR CYLINDER

Electromagnetic scattering from an infinite cylinder is a two dimensional problem. Its solution has been established for decades. For finite cylinders, Storer [Ref. 5] deals with the case of long thin cylinder (a wire) and Kennedy [Ref. 6] with a short thick cylinder (a disc). In practical applications the finite structures are the cases of interest.

The geometric diffraction theory has been used by Ross [Ref. 7] to calculate the EM scattering from a finite cylinder where the cylinder was 25λ long and 5λ in diameter. This approach can only deal with problems near the optical region. When dealing with problems in the resonance region the complete Maxwell equations have to be used. One approach to setup the boundary value problem is by using the Chu-Stratton integral to formulate an integrodifferential equation.

The solution that will be shown in this chapter is for circular tubular cylinder having very thin walls. Cylinder with those specifications were tested in the Scattering Laboratory at the Naval Postgraduate School and the measurement data as shown in Chapter III was used to compare with this theory.

The finite cylinder is assumed to be an infinitesimally thin walled, perfectly conducting circular tube of radius "a" and length "2h". Its center is located at the origin of the cartesian coordinates (x,y,z); its axis coincides with the z axis of the coordinates, so that the length extends from z=-h to z=+h; Figure 2.2 . To simplify the equations, the scaled cylindrical coordinates (ρ,ϕ,z) have been used (Figure 2.3) in which the cylinder length extends from z=-l to z=+1 and its radius is 1.

The surface current can flow only on the perfectly conducting surfaces when the free space case is assumed; and because of the thin walls the total surface current can be assumed to exist only for $-1 \le z \le +1$ and $\rho=1$. There are only two current components: the circumferential current circles around the cylinder in the ϕ direction while the axial current travels along the cylinder in the z direction. Both currents are functions of the position on the cylinder. In terms of their Fourier coefficients the axial current density $K_z(\phi,z)$ and the circumferential current density $K_{\varphi}(\phi,z)$ can be represented as equations 2.2 and 2.3 .

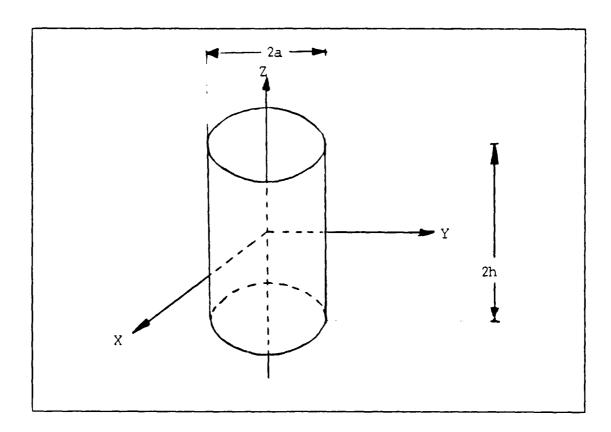


Figure 2.2 Cylinder in Cartesian Coordinates

$$K_{z}(\phi,z) = \sum_{n=-\infty}^{\infty} K_{zn}(z) \exp\{in\phi\} =$$
 (2.2)

$$= \sum_{n=0}^{\infty} K_{zn}^{(+)}(z) \cos(n\phi) + iK_{zn}^{(+)}z) \sin(n\phi)$$

Where-
$$K_{z0}^{(+)}(z)=K_{z0}(z)$$
 $K_{z0}^{(-)}(z)=0$ (2.2 a)

$$K_{zn}^{(\pm)}(z) = K_{zn}(z) \pm K_{z(-n)}(z)$$
 $n \neq 0$ (2.2 b)

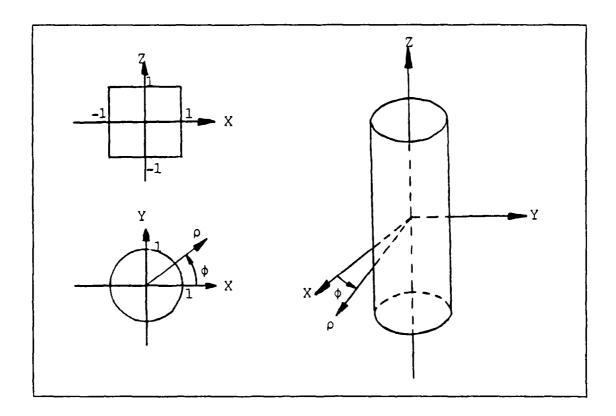


Figure 2.3 Cylinder in Cylindrical Coordinates

$$K_{\phi}(\phi,z) = \sum_{n=-\infty}^{\infty} K_{\phi n}(z) \exp\{in\phi\} =$$

$$= \sum_{n=0}^{\infty} K_{\phi n}^{(*)}(z) \cos(n\phi) + iK_{\phi n}^{(-)}(z) \sin(n\phi)$$
here-
$$K_{\phi 0}^{(+)}(z) = K_{\phi 0}(z) \qquad K_{\phi 0}^{(-)}(z) = 0 \qquad (2.3 a)$$

Where-

$$K_{\phi n}^{(\pm)}(z) = K_{\phi n}(z) \pm K_{\phi (-n)}(z)$$
 $n \neq 0$ (2.3 b)

On the surface of the cylinder z takes on values in the range of $-1 \le z \le +1$ and thus we can represent z by $z = \cos(v)$ where $0 \le v \le \pi$. Near the edges $z = \pm 1$, the current density in ϕ direction approaches $(1-z^2)^{-1/2}$ and in z direction approaches $(1-z^2)^{+1/2}$ because of the edge conditions. This leads to the definitions of K_{zn}^p and $K_{\phi n}^p$ in equations 2.4 and 2.5 . $K_{zn}^{(\pm)p}$ can be defined by 2.2 a ,2.2 b; and $K_{\phi n}^{(\pm)p}$ by 2.3 a ,2.3 b.

$$K_{zn}(z) = (1/\pi) \sum_{p=0}^{\infty} K_{zn}^{p} \sin[(p+1)v]$$
 (2.4)

$$K_{\phi n}(z) = (1/\pi \sin v) \sum_{p=0}^{\infty} K_{\phi n}^{p} \cos(pv)$$
 (2.5)

The surface current that described here in terms of $K_{2n}^{\ p}$ 2.4 and $K_{\phi n}^{\ p}$ in 2.5 is the sum of the inside and outside surface currents. The inside surface current is on $\rho=1^+$ while the outside surface current is on $\rho=1^+$. The radiations due to these currents together with the incident electric field E_i should satisfy the boundary conditions of Maxwell's theory. For a tubular cylinder in a medium with homogeneous, isotropic permittivity ϵ and permeability μ the relations among the tangential components of the electric fields and the surface current can be written as 2.6 , 2.7 , 2.8 and 2.9 [Ref. 8].

$$(1 + \frac{1}{1_{1}^{2}} \frac{\lambda^{2}}{\partial z^{2}}) \int_{-1}^{1} dz_{0} K_{zn}(z_{0}) G_{n}(1_{1} | z - z_{0} |, 1_{2})$$
 (2.6)

$$+\frac{in}{1_1 \cdot 1_2} \frac{\partial}{\partial z} \frac{1}{1_1} dz_0 K_{\phi n}(z_0) G_n(1_1 | z - z_0 | , 1_2) =$$

$$=-(2i/1_1 1_2 \zeta_0) E_{zn}^{s}(z)$$

$$\frac{1}{1} dz_{0} K_{\phi n}(z_{0}) \{1/2[G_{n-1}(1_{1}|z-z_{0}|,1_{2}) + G_{n+1}(1_{1}|z-z_{0}|,1_{2})] - \frac{n^{2}}{12} G_{n}(1_{1}|z-z_{0}|,1_{2})\}
+ (in/1_{1}1_{2}) \frac{\partial}{\partial z_{-}} \int_{1}^{1} dz_{0} K_{zn}(z_{0}) G_{n}(1_{1}|z-z_{0}|,1_{2})
= -(2i/1_{1}1_{2}\zeta_{0}) E_{\phi n}^{s}(z)$$

$$E_{zn}^{s}(z) + E_{zn}^{i}(z) = 0 \qquad -1 < z < +1 \qquad (2.8)$$

Where: $l_1=kh=2\pi h/\lambda$

$$l_2 = ka = 2\pi a/\lambda$$

$$\zeta_0 = (\mu/\epsilon)^{1/2}$$

$$G_n(1_1|z-z_0|,1_2)=$$

$$\int_{-\pi}^{\pi} (d\phi/2\pi) \exp[-in(\phi-\phi_0)]G[l_1|z-z_0|,2l_2|\sin(\phi-\phi_0)/2|]$$

$$G(x_1,x_2) = {\exp[i(x_1^2 + x_2^2)^{1/2}]}/(x_1^2 + x_2^2)^{1/2}$$

Equations 2.6 and 2.7 can be obtained from the Stratton-Chu equations [Refs. 9,4 pp. 99-107,464], together with the edge condition that $K_z(\phi,z)=0(1-z^2)^{1/2}$ as $|z| \rightarrow 1$. Equations 2.8 and 2.9 are boundary conditions for the tangential electric field components on a perfectly conducting surface.

Equation 2.6 , 2.7 , 2.8 and 2.9 give the connection between the current density on the cylinder surface and the electric fields on the surface. For the back scattering cross section, the far field should be obtained. Denote an arbitrary point in the far field (ρ,ϕ,z) by the vector r and a point on the cylinder surface by the vector r_0 , then in the far field:

$$k|\vec{\tau}| = (1_2^2 \rho^2 + 1_1^2 z^2)^{1/2} >> 1$$

$$k|\vec{\tau}| >> (1_1^2 + 1_2^2)^{1/2} \ge k|\vec{\tau}_0|$$
(2.10)

The scattered far field at point r will be 2.11 , 2.12 and 2.13 where θ is the angle between the z axis and the vector r.

$$-(2i/1_{1}1_{2}z_{0})E_{z}(\rho,\phi,z) =$$

$$= \sin^{2}\theta \int_{-1}^{1} dz_{0} \int_{-\pi}^{\pi} (d\phi_{0}/2\pi)G(\vec{r}-\vec{r}_{0})K_{z}(\phi_{0},z_{0})$$

$$-\sin^{2}\cos^{2}\theta \int_{-1}^{1} dz_{0} \int_{-\pi}^{\pi} (d\phi_{0}/2\pi)G(\vec{r}-\vec{r}_{0})\sin(\phi-\phi_{0})K_{\phi}(\phi_{0},z_{0})$$

$$-(2i/1_11_2\zeta_0)E_{\rho}(\rho,\phi,z) = (2.12)$$

=-
$$\sin\theta\cos\theta \int_{-1}^{1} dz_0 \int_{-\pi}^{\pi} (d\phi_0/2\pi) G(\vec{r} - \vec{r}_0) K_z(\phi_0, z_0)$$

+cos²
$$\theta$$
 $\int_{-1}^{1} dz_0 \int_{-\pi}^{\pi} (d\phi_0/2\pi)G(\vec{\tau}-\vec{\tau}_0)\sin(\phi-\phi_0)K_{\phi}(\phi_0,z_0)$

$$-(2i/1_11_2\zeta_0)E_{\phi}(\rho,\phi,z)= \qquad (2.13)$$

$$= \int_{-1}^{1} dz_{0} \int_{-\pi}^{\pi} (d\phi_{0}/2\pi)G(\vec{r} - \vec{r}_{0}) \cos(\phi - \phi_{0})K_{\phi}(\phi_{0}, z_{0})$$

To simplify the equations, spherical coordinate would be used; with the electric field components $E_r(r,\theta,\varphi)$, $E_\theta(r,\theta,\varphi)$ and $E_\varphi(r,\theta,\varphi)$.

Since:

$$E_{r}(r,\theta,\phi)=E_{\rho}(\rho,\phi,z)\sin\theta+E_{z}(\rho,\phi,z)\cos\theta \qquad (2.14)$$

$$E_{\theta}(r,\theta,\phi) = E_{\rho}(\rho,\phi,z)\cos\theta - E_{z}(\rho,\phi,z)\sin\theta \qquad (2.15)$$

The field components in the spherical coordinate are:

$$(-2i/1_1 1_2 \zeta_0) E_r(r, \theta, \phi) = 0$$
 (2.16)

$$(-2i/1_11_2\zeta_0)E_{\theta}(r,\theta,\phi)=$$
 (2.17)

$$=-\sin\theta\int_{-1}^{1}dz_{0}\int_{-\pi}^{\pi}(d\phi_{0}/2\pi)G(\vec{r}-\vec{r}_{0})K_{z}(\phi_{0},z_{0})$$

$$+\cos\theta \int_{-1}^{1} dz_{0} \int_{-\pi}^{\pi} (d\phi_{0}/2\pi)G(\vec{r}-\vec{r}_{0})\sin(\phi-\phi_{0})K_{\phi}(\phi_{0},z_{0})$$

$$(-2i/1_11_2\zeta_0)E_{\phi}(r,\theta,\phi)=$$
 (2.18)

$$= \int_{-1}^{1} dz_{0} \int_{-\pi}^{\pi} (d\phi_{0}/2\pi)G(\vec{r} - \vec{r}_{0})\cos(\phi - \phi_{0})K_{\phi}(\phi_{0}, z_{0})$$

Where $G(\vec{r}-\vec{r}_0)$ are approximated by 2.19 in the far field.

$$G(\vec{r} - \vec{r}_0) = \qquad (2.19)$$

$$= G(\vec{r}) \exp[-il_2 \sin\theta \cos(\phi - \phi_0)] \exp[-il_1 \cos\theta z_0]$$

$$G(\vec{r}) = \exp[ikr]/kr$$

Since:

$$\int_{-\pi}^{\pi} (d\phi/2\pi) \cos(n\phi) \exp[-il_2 \sin\theta \cos(\phi-\phi_0)] = i^{-n} J_n(l_2 \sin\theta)$$

and

$$\int_{-1}^{1} (dz/\pi \sqrt{1-z^{2}}) \cos(pv) \exp[-il_{1}\cos(\theta)z_{0}] =$$

$$= \int_{0}^{\pi} (dv/\pi) \cos(pv) \exp[il_{1}\cos(\theta)\cos v] = i^{-p} J_{p}(l_{1}\cos\theta)$$

and by using K_{2n}^p and $K_{\varphi n}^p$ from equations 2.4 and 2.5 , the fields in the far field region can be written as equations 2.20 2.21 and 2.22 .

$$\mathbf{E}_{\mathbf{r}}(\mathbf{r},\theta,\phi)=0\tag{2.20}$$

$$[-2i/1_1 1_2 \zeta_0 G(\vec{r})] E_{\theta}(r, \theta, \phi) =$$
 (2.21)

$$=-\sum_{n=0}^{\infty}\sum_{p=0}^{\infty}i^{-(n+p)}\left[(p+1)\sin\theta/1_{1}\cos\theta\right]J_{p+1}(1_{1}\cos\theta)J_{n}(1_{2}\sin\theta)$$

$$[K_{zn}^{(+)} {\overset{p}{\sim}} cos(n_{\phi}) + iK_{zn}^{(-)} {\overset{p}{\sim}} sin(n_{\phi})]$$

$$\begin{array}{ll} + \sum\limits_{n=1}^{\infty} \sum\limits_{p=0}^{\infty} i^{-(n+p)} [n cos\theta/l_2 sin\theta] J_p(l_1 cos\theta) J_n(l_2 sin\theta) \end{array}$$

$$[K_{\phi n}^{(-)} \mathcal{P}_{cos}(n\phi) + iK_{\phi n}^{(+)} \mathcal{P}_{sin}(n\phi)]$$

$$[-2i/1_11_2\zeta_0G(\vec{r})]E_{\phi}(r,\theta,\phi)=$$
 (2.22)

$$= \sum_{n=0}^{\infty} \sum_{p=0}^{\infty} (n \cdot p) J_p(1_1 \cos \theta) J_n'(1_2 \sin \theta) [-K_{\phi n}^{(-)p} \sin(n_{\phi}) + iK_{\phi n}^{(+)p} \cos(n_{\phi})]$$

At this point the scattered field in the far field region is written in terms of the Fourier series expansion of the surface current flowing on the cylinder. To calculate the back scattering cross section of a cylinder the incident wave should be inspected as in equation 2.1 . With a linearly polarized plane incident wave on the target having unit strength and zero phase at the center of the target, the cross section is given by squation 2.23 and the phase shift is given by squation 2.24 .

$$\sigma = \lim_{r \to \infty} 4\pi \, r^2 \, |E_s|^2 \tag{2.23}$$

$$\delta = \arg\{E_{s} \exp[-ikr]\}$$
 (2.24)

In the special case of broadside incidence with polarization along the cylinder axis, which contained the z axis, the propagation vector k is given by equation 2.25 and the incident field on the surface of the cylinder by equation 2.26.

$$\hat{\mathbf{k}} = \mathbf{k}\hat{\mathbf{x}} \tag{2.25}$$

$$\vec{E}_{i} = \hat{z} \cdot \exp[ikx] = \hat{z} \cdot \exp[ika\cos(\phi)]$$
 (2.26)

Since E_z^i is an even function in ϕ , $K_{zn}^{(-)}(z)=0$ and from 2.6 and 2.7, only $K_{\phi n}^{(-)}(z)$ is coupled to $K_{zn}^{(+)}(z)$. Thus $K_{\phi n}^{(+)}=0$.

In the far field the back scattered field has the components E_θ and E_φ as shown on 2.27 and 2.28 and a total back scattered field as in 2.29 .

$$[-2i/1_1 1_2 \zeta_0 G(\vec{r})] E_{\theta}^{s}(r,\theta,\phi) =$$
 (2.27)

$$=\sum_{n=0}^{\infty}\sum_{p=0}^{\infty}i^{-n}(-1)^{p+1}[(2p+1)\sin\theta/1_{1}\cos\theta]$$

$$J_{2p+1}(1_1\cos^{\theta})J_n(1_2\sin^{\theta})K_{2n}^{(+),2p}\cos(n\phi) + \sum_{n=1}^{\infty}\sum_{p=0}^{\infty}i^{-n+1}(-1)^{p+1}$$

$$[\mathsf{ncos}_\theta/\mathsf{1}_2\mathsf{sin}_\theta)\mathsf{J}_{2p+1}(\mathsf{1}_1\mathsf{cos}_\theta)\mathsf{J}_n(\mathsf{1}_2\mathsf{sin}_\theta)\mathsf{K}_{\phi n}^{(-)\,2p+1}\;\mathsf{cos}(n_\phi)$$

$$[-2i/1_1 1_2 \zeta_0 G(\vec{r})] E_{\phi}^{s}(r,\theta,\phi) =$$
 (2.28)

$$= \sum_{n=1}^{\infty} \sum_{p=0}^{\infty} i^{-n+1} (-1)^p J_{2p+1}(1_1 \cos \theta) J_n'(1_2 \sin \theta) K_{\phi n}^{(-)2p+1} \sin(n\phi)$$

$$[-2i/1_{1}1_{2}\zeta_{0}G(\hat{r})]\hat{z}\vec{E}_{sc}(r,\pi/2,\pi) =$$

$$=+1/2\sum_{n=0}^{\infty}i^{n}J_{n}(1_{2})K_{zn}^{(+),0}$$
(2.29)

Using equations 2.23 and 2.24 with 2.29, the back scattered cross section and the phase shift of a tubular cylinder is given by 2.30 and 2.31.

$$\sigma/ah = (2.30)$$

$$= \lim_{r \to \infty} (4\pi r^{2}/ah) \left| \left[1_{1} 1_{2} \zeta_{0} G(\vec{r}) / -2i \right] 1 / 2 \sum_{n=0}^{\infty} i^{n} J_{n} (1_{2}) K_{zn}^{(+)}, 0 \right|^{2}$$

$$= (4\pi r^{2}/ahk^{2}r^{2}) \left| (1_{1} 1_{2} \zeta_{0} / 4) \sum_{n=0}^{\infty} i^{n} J_{n} (1_{2}) K_{zn}^{(+)}, 0 \right|^{2}$$

$$= \left| (4\pi^{2} \zeta_{0} / 1_{1} 1_{2}) \sum_{n=0}^{\infty} i^{n} J_{n} (1_{2}) K_{zn}^{(+)}, 0 \right|^{2}$$

$$\delta = (2.31)$$

$$= \arg\{ [\exp(-ikr)G(\vec{r}) 1_1 1_2 \zeta_0 / -2i] 1/2 \sum_{n=0}^{\infty} i^n J_n (1_2) K_{2n}^{(+),0} \}$$

$$= \arg\{ [\exp(-ikr)G(\vec{r}) / i] \sum_{n=0}^{\infty} i^n J_n (1_2) K_{2n}^{(+),0} \}$$

$$= \arg[\sum_{n=0}^{\infty} i^{n-1} J_n (1_2) K_{2n}^{(+),0}]$$

From equation 2.30 one can see that the back scattered cross section of a tubular cylinder at the broadside aspect angle is depended only on the Fourier components of the axial surface current along the z direction. That assumption was tested against the experimental results and shown in Chapter IV.

III. MEASUREMENTS

The exact solutions to the back scattering cross sections of targets are known only for few bodies. For almost all cases approximate solutions have to be sought. Experimental verification is the only justification for a good approximation.

The back scattering measurement facility in the Scattering Laboratory of the Naval Postgraduate School is an indoor range designed for model measurements above 2GHz. The distance from the antennas to the target, the target itself and the radar output wavelength are scaled down from life size. One advantage of using the indoor range is the practicability of testing models that are smaller and cheaper than full scale targets, even though tighter specifications on target details have to be met. This chapter deals with the experimental setup, the targets and the measurement procedures.

A. SET-UP

Figure 3.1 is a block diagram showing the signal flow of the setup. Table 1 is a list of the equipment in the setup.

The frequency and output power level of the signal generator HP-8672A is controlled by an HP-85 Microcomputer. The output RF signal from the signal generator enters an GaAs wide band amplifier which is operated at saturation to provide an output of about 23dBm. The amplified RF signal is then passed through the directional coupler to feed the transmitting horn antenna. The returned signal from the scattered field of the target is collecter by the receiving horn antenna to feed the test port of the harmonic frequency

TABLE 1
Equipment

| Microcomputer | HP-85 |
|------------------------------|-------------------|
| Flexible Disk Drive | HP-82901M |
| Plotter . | HP-7225B |
| Synthesized Signal Generator | HP-8672A |
| RF Amplifier | Avanter SA83-2953 |
| Dual DC Power Supply | HP-6227B |
| Digital Multimeter | HP-3466A |
| Directional Coupler | Narda 5292 |
| Transmitting Antenna | |
| Receiving Antenna | |
| Harmonic Frequency Converter | HP-8411A |
| Network Analyzer | HP-8410C |
| Phase Magnitude Display | HP-8412B |
| Digital Voltmeter | HP-3455A |
| Digital Voltmeter | HP-3456A |
| HP Interface Bus | |
| Coaxial Cables | |

converter, HP-8411A. The portion of the transmitting signal coupled out through the directional coupler, is attenuated to get 43dB attenuation before it is fed to the reference port of the harmonic frequency converter. The harmonic frequency converter and the network analyzer HP-8410C, with the phase-magnitude display HP-8412B function as a phase difference and magnitude ratio meter between the transmitting and receiving signals. The phase difference and the magnitude ratio are converted to volts that are measured by the digital voltmeters HP-3456A and HP-3455A. The digital word from the voltmeters is then transferred to the microcomputer for processing and then displayed on the plotter HP-82901M.

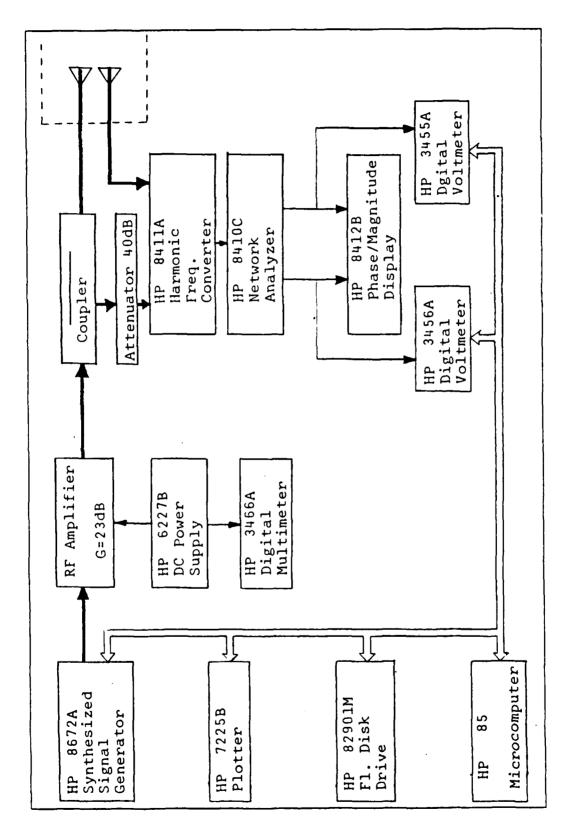


Figure 3.1 Signal Flowing Diagram

The antennas used in this system are identical linearly polarized horn antennas. The scattered electric field from a target has three components. An antenna will pick a linear combination of the components and couple it to the receiver through a transmission line.

The polarization of the transmitting and receiving antennas can be represent by the matrixes shown in equations 3.1 and 3.2

$$\hat{q} = \begin{bmatrix} \cos \gamma_t \\ \sin(\gamma_t) \exp(i\delta_t) \end{bmatrix}$$
 (3.1)

$$\hat{p} = [\cos \gamma_r, \sin \gamma_r \exp(i \delta_r)]$$
 (3.2)

Where \hat{q} -Unit column matrix defining the polarization of transmitting antenna

p-Unit row matrix defining the polarization of receiving antenna

γ-An angle which donates the orientation of the linear polarization referred to the horizontal plane

δ-Phase angle

t-Denotes transmitting antenna

r-Denotes receiving antenna

For linear horizontal polarization the matrices are given by equation 3.3 and 3.4

$$\hat{\mathbf{q}} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \tag{3.3}$$

$$\hat{p} = [1,0]$$
 (3.4)

The target has a complex scattering matrix which is a function of the geometry of the target and the frequency of the incident wave. Assuming a uniform plane incident wave at the target, the scattering matrix is a linear relation between the incident electric field and the scattered far field from the target. The Maxwell equations are linear as long as ε and μ are linear. This is true even if ε and μ are unisotropic and inhomogeneous. The scattering matrix can be written as 3.5

$$S = \begin{bmatrix} \sqrt{\sigma_{HH}} \exp(i\rho_{HH}) & \sqrt{\sigma_{HV}} \exp(i\rho_{HV}) \\ \sqrt{\sigma_{VH}} \exp(i\rho_{VH}) & \sqrt{\sigma_{VV}} \exp(i\rho_{VV}) \end{bmatrix}$$
(3.5)

Where $\sqrt{\sigma}$ -Magnitude of the scattering matrix element ρ -Phase of the scattering matrix element H-Denotes horizontal polarization v-Denotes vertical polarization

The radar cross section of a target with scattering matrix S and obtained by a pair of transmitting and receiving antennas with polarization \hat{q} and \hat{p} respectively will be 3.6

$$\sigma = |\hat{p}\hat{Sq}|^2 \tag{3.6}$$

In our system, the antennas are horizontally polarized so that:

$$\sigma = \sigma_{HH} \tag{3.7}$$

The characteristics of these antennas are given in Appendix A. These antennas are mounted on the wall of an anechoic chamber in the Scattering Laboratory at the Naval Postgraduate School. All measurements were taken inside the chamber. The characteristics of the anechoic chamber is given in Appendix B and it was discussed in great detail by Mariategui [Ref. 10].

B. TARGETS

Figure 3.2 shows the orientation of the cylinder in the anechoic chamber. The antennas-to-target distance was two meters. The targets were a set of tubular circular cylinders made of thin walled brass of various lengths and diameters. Figure 3.3 shows the dimensions of the cylinders.

The back scattering cross section of a tubular cylinder depends on the following parameters:

- (1) The cylinder length (2h).
- (2) The cylinder diameter (2a).
- (3) The cylinder wall thickness.
- (4) Azimuth aspect angle.
- (5)Cylinder tilt angle.
- (6) Transmitting antenna polarization.
- (7) Receiving antenna polarization.
- (8) Transmitter frequency (f).

In the measurements the effects of varying wall thickness was neglected because it was small compared to the wavelengths and other dimensions of the cylinders. The polarization was always parallel to the axis of the cylinder. The remaining three parameters were varied and their effects on the back-scattering cross section of the cylinder were studied. Table 2 gives the characteristics of the targets used in this experiment.

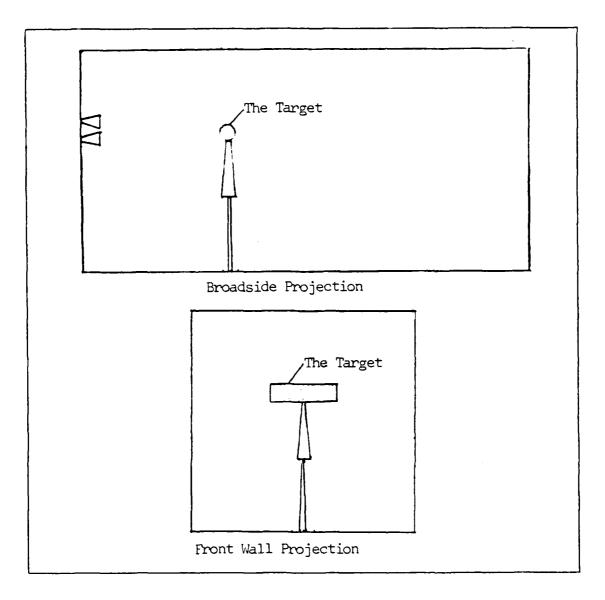


Figure 3.2 Orientation of Targets in Anechoic Chamber

Dimensions of TARGET20 and TARGET21 are shown in Figure 3.4. These targets are cylinders with four rectangular fins having the dimensions: 0.75x0.375x0.01 inches. The axis of the cylinder coincides with the z axis in both targets while the fins are on the x-y axiss in TARGET20 and 45 degrees of the axiss in TARGET21.

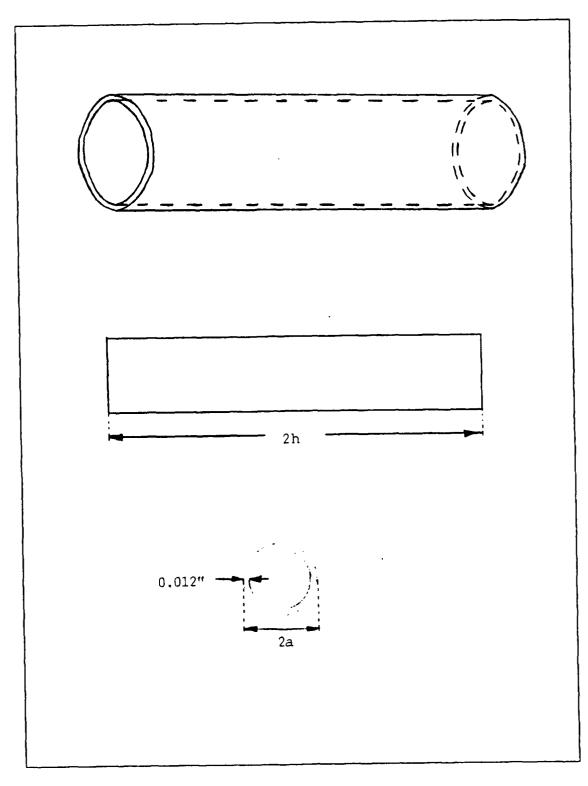


Figure 3.3 Cylinders Dimensions

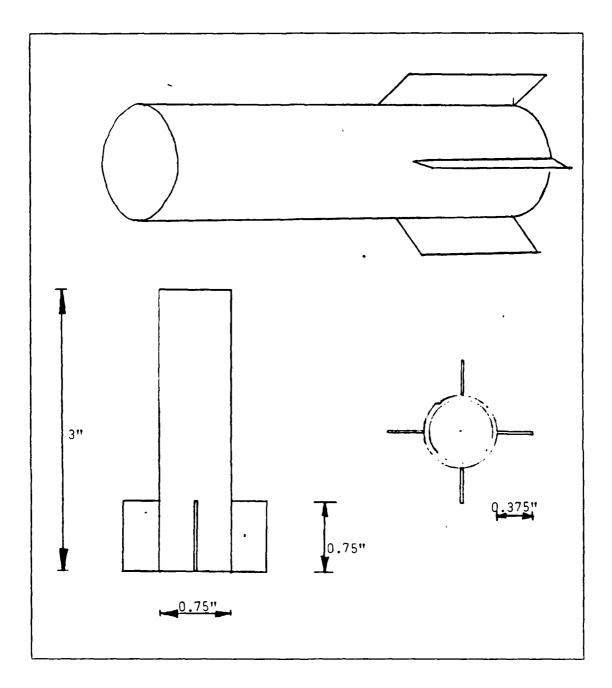


Figure 3.4 The Fins Orientation

TABLE 2
Targets Description

| Name | Description | Length (2h) | Diameter (2a) |
|----------|--------------------|-------------|---------------|
| | | in inches | in inches |
| TARGET1 | Plane cylinder | 2.0 | 0.375 |
| TARGET2 | Plane cylinder | 2.0 | 0.5 |
| TARGET3 | Plane cylinder | 2.0 | 0.75 |
| TARGET4 | Plane cylinder | 2.25 | 0.375 |
| TARGET5 | Plane cylinder | 2.25 | 0.5 |
| TARGET6 | Plane cylinder | 2.25 | 0.75 |
| TARGET7 | Plane cylinder | 2.5 | 0.375 |
| TARGET8 | Plane cylinder | 2.5 | 0.5 |
| TARGET9 | Plane cylinder | 2.5 | 0.75 |
| TARGET10 | Plane cylinder | 2.75 | 0.375 |
| TARGET11 | Plane cylinder | 2.75 | 0.5 |
| TARGET12 | Plane cylinder | 2.75 | 0.75 |
| TARGET13 | Plane cylinder | 3.0 | 0.375 |
| TARGET14 | Plane cylinder | 3.0 | 0.5 |
| TARGET15 | Plane cylinder | 3.0 | 0.75 |
| TARGET16 | Plane cylinder | 1.5 | 0.375 |
| TARGET17 | Plane cylinder | 4.5 | 0.75 |
| TARGET18 | Plane cylinder | 2.5 | 0.625 |
| TARGET19 | Plane cylinder | 3.75 | 0.625 |
| TARGET20 | Cylinder with fins | 3.0 | 0.75 |
| TARGET21 | Cylinder with fins | 3.0 | 0.75 |

C. MEASUREMENT PROCEDURE

The cross section is defined in terms of a uniform plane incident wave. The radiation of the antenna approximates that of a dipole and is not a plane wave. To get a good plane wave approximation, a distance that satisfied 3.8, 3.9 and 3.10 was chosen. The relationship shown in 3.10 assumed difference of no more then 20 degrees between the phase at the center and the edges of the target [Ref. 11].

$$r \ge 10\lambda \tag{3.8}$$

$$r \ge 2D^2/\lambda \tag{3.10}$$

Where r-Antenna to target distance.

D-Maximum dimension of the target.

 λ -Wavelength of the transmitted wave.

The signal picked up by the receiving antenna is the vectoral sum of the target echo and the background radiation. To cancel the coupling between the antennas and the direct back scattering from the support and the walls, background measurement was carried out by measuring the echo returned when the target was not present. The difference between signals when the target was present and when the target was absent gave the echo signal of the target.

Calibration of the system to take into account the characteristics of the system which are dependent on frequencies and to relate the echo signal power to the target back scattering cross section and phase shift was achieved by

measuring the echo signal from a sphere and comparing the experimental results to theoretical values. The theoretical data was calculated by Mie series computed with the Sphere program [Appendix C]. The calibration was done with the Calib program [Appendix D]. Measurements of back scattering of the target were done with the Target program [Appendix E]. Description and explanations of the computer programs were given by Lolic [Ref. 12].

The instruments add noise to the measurement data. This noise is white noise and to minimize its effects, each target was measured five times and averaged. Before each measurement new calibration was done to minimize the noise effects in the calibration. For phase shift near ±180 degrees the average procedure did not take care of the discontinuity properly and the averaged result depended on the number of times the measured values took on +180 or -180 degrees.

Measurement errors that could not be controled are the coupling between the target and the support and the bistatic coupling which is the strong target scattered lobs in the forward hemisphere (away from the antennas) and then back to the receiving antenna from the walls. The latter effect is believed to have been taken care of through the use of the microwave absorbers.

D. MEASURED DATA

The measured results of the 21 targets were plotted and given at the end of this chapter (Figure 3.5 to 3.46). Those plots contained cross section and phase shift versus frequency data, for each target. The theoretical data are given in Tables 3 to 23. The data was used for comparison with theoretical values and it shown in Chapter IV.

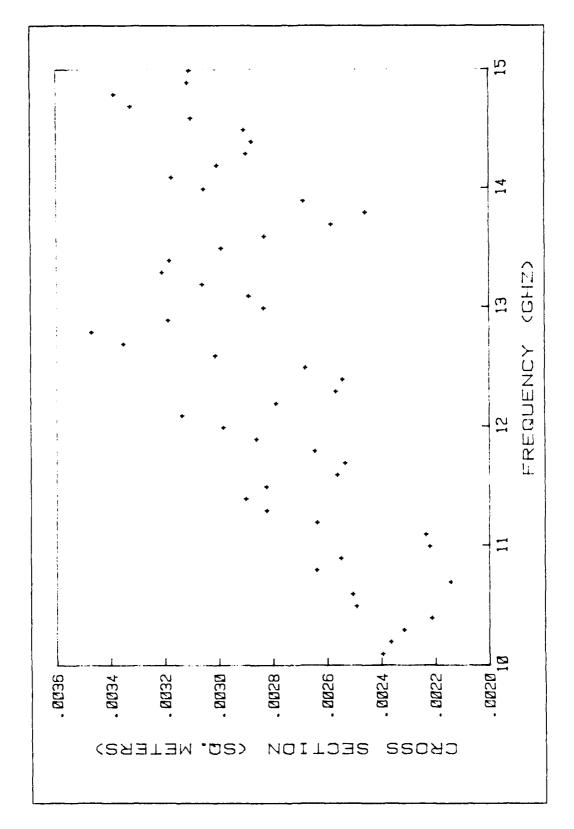


Figure 3.5 TARGET1 Cross-Section vs. Frequency

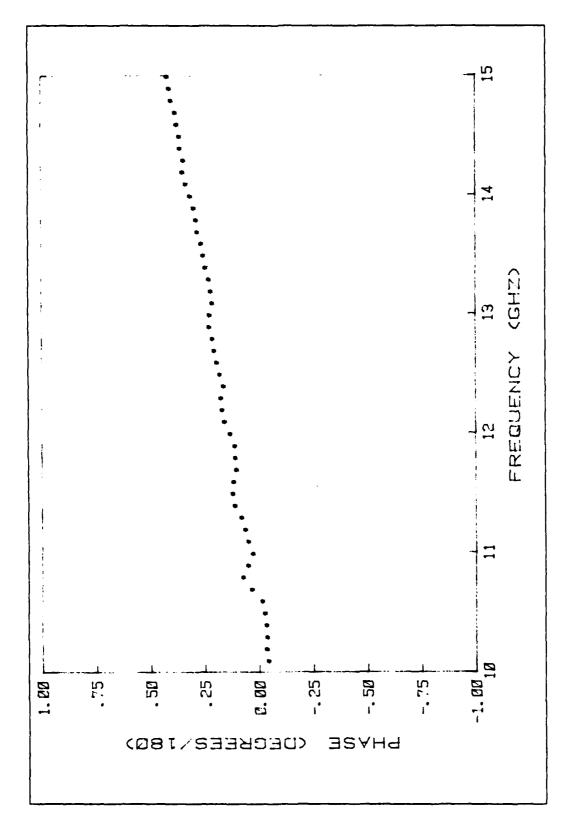


Figure 3.6 TARGET1 Phase Shift vs. Frequency

TABLE 3
TARGET1 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|--|--------------------------|--------------------------|
| 10.10 10.20 | .00238 00235 | - 05747 - 04705 |
| 10 30 | 00235 .00230 | - 05245 |
| 10.40 10.50 | .00220 .00248 | - 04795 - 04077 |
| 10.60 | .00249 | - 02920 |
| 10 70 10.80 | 00213 00263 | .01917 05879 |
| 10.90 | 99254 | 03643 |
| 11.00 11.10 | .00221 .00222 | .01517 03640 |
| 11.20 | 00262 | .04978 .06599 |
| 11 30 11.40 | .00281 .00289 | 09669 |
| 11.50 | 00281 | .19508 .19195 |
| 11.60 11.70 | . 99255 . 99252 | 08954 |
| 11.80 | 00263 00265 | .09357 .09612 |
| 11 90 12 00 | .00297 | .11657 |
| 12.00 12.10 12.20 | 00312 | .14314 .15434 |
| 12.20 | .00278 00255 | 15929 |
| 12.20 12.30 12.40 12.50 12.60 12.70 | .00253 .00267 | .14789 .16398 |
| 12.50 12.60 | 90300 | . 17882 |
| 12.70 | .00334 .00346 | .18917 19790 |
| | 00317 | 19790 21208 |
| 13.00 13.10 13.20 13.30 13.40 | 00282 00288 | .20998 .19925 |
| 13.10 | . 00305 | . 20384 |
| 13.30 | 00320 00317 | 21406 22914 |
| 13 50 | .00298 | 24056 |
| 47 66 | .00282 .00257 | 25054 26677 |
| 13.70 13.80 13.90 | .00244 | 27079 28132 |
| 13.90 14.00 | .00267 .00304 | 29834 |
| 14 10 | 00316 | 31943 33389 |
| 14.20 14.30 | .00299 00288 | .33000 |
| 14.40 | .00286 | .34395 .34727 |
| 14 50 14.60 | .00289 00309 | . 35859 |
| 14.70 | .00331 | .36647 .386 51 |
| 14.80 14.90 | 00337 .00310 | . 39357 |
| 15.00 | 00309 | 40518 |

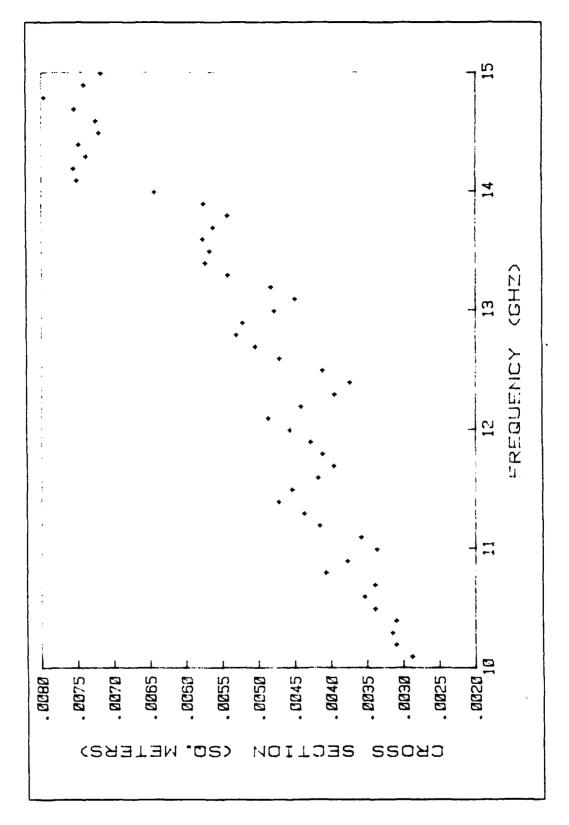


Figure 3.7 TARGET2 Cross-Section vs. Frequency

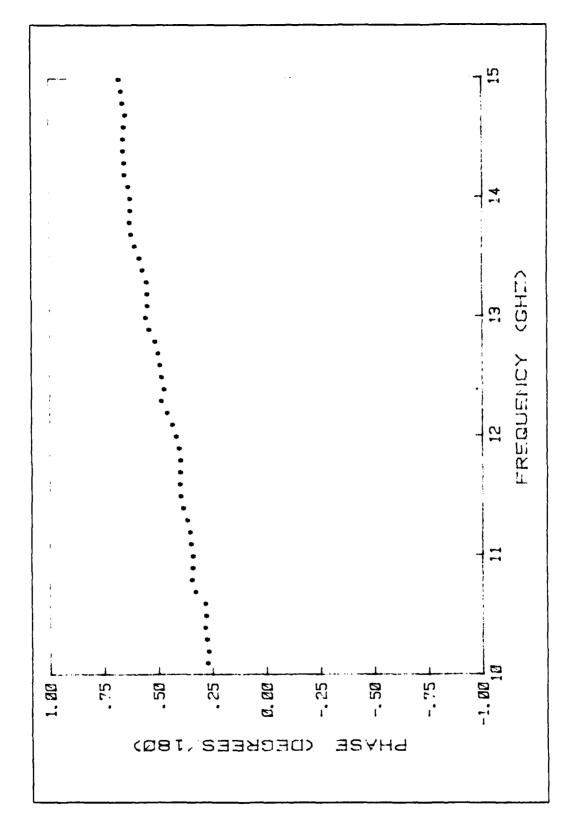


Figure 3.8 TARGET2 Phase Shift vs. Frequency

TABLE 4
TARGET2 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|---|---|--|
| 100 340 000 000 000 000 000 000 000 000 0 | 0023334940233413793350000000000000000000000000000000000 | 914140301899987443319552921113480956638684691586 22222237968869744487733468495663386844691586 222223746813331944467733448896363333444691586 2333333333333333333333333333333333333 |

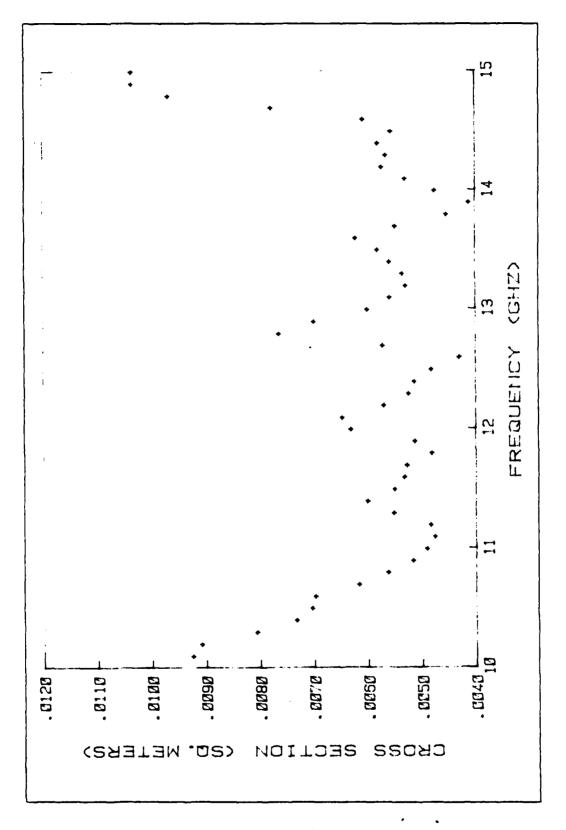


Figure 3.9 TARGET3 Cross-Section vs. Frequency

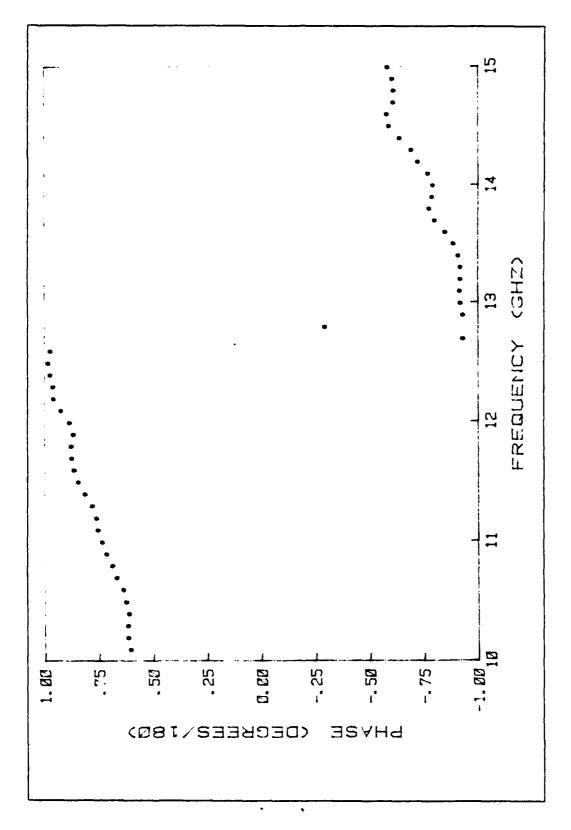


Figure 3.10 TARGET3 Phase Shift vs. Frequency

TABLE 5
TARGET3 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|---|--|---|
| 100.000.000.000.000.000.000.000.000.000 | 913 913 913 913 913 913 913 913 913 913 | 16611410734705639524773755653031689,4604765152508531691334705639524773755653031689,460476515250853165910233470956530944619935227297644210537377352485316657624476031602330044619645320333312442193737735438266577776692235399999999988888888899999999988888877776692219 |

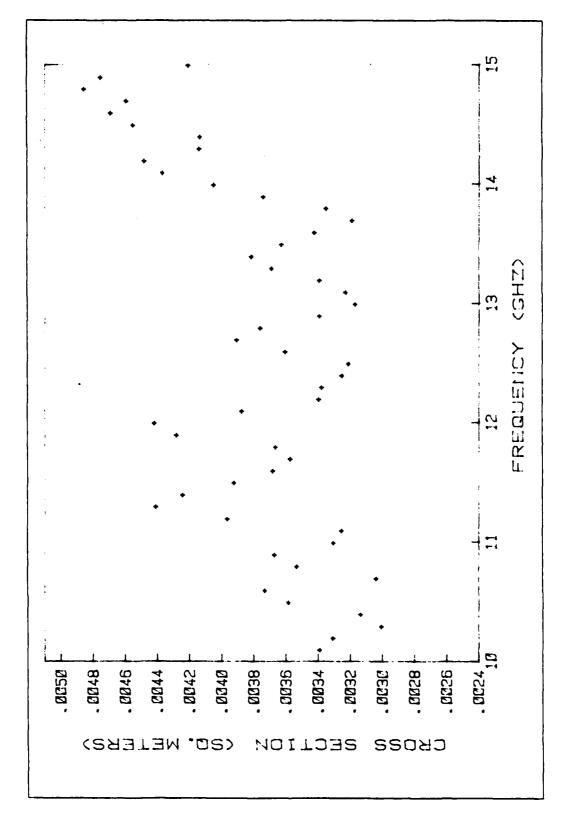


Figure 3.11 TARGET4 Cross-Section vs. Frequency

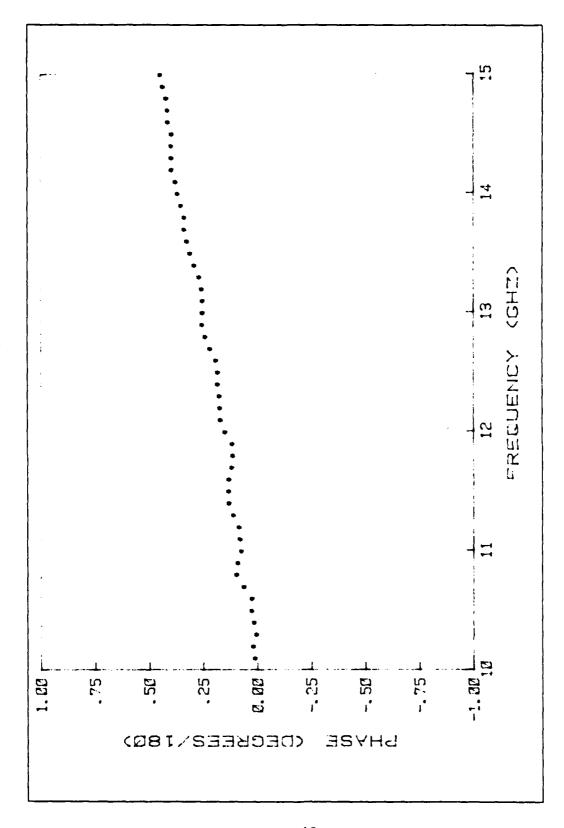


Figure 3.12 TARGET4 Phase Shift vs. Frequency

TABLE 6
TARGET4 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|--|---|----------------------|
| 10000000000000000000000000000000000000 | 799271215934920655606863332199947751177010732356227373235622737323562273732356227373235622737323562273732356227373235622737323562273732356222778444444444444444444444444444444444 | |

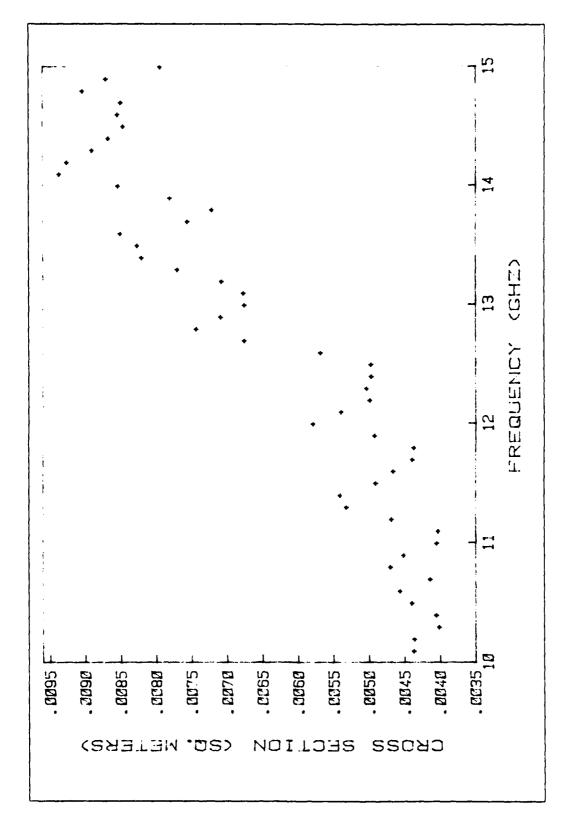


Figure 3.13 TARGET5 Cross-Section vs. Frequency

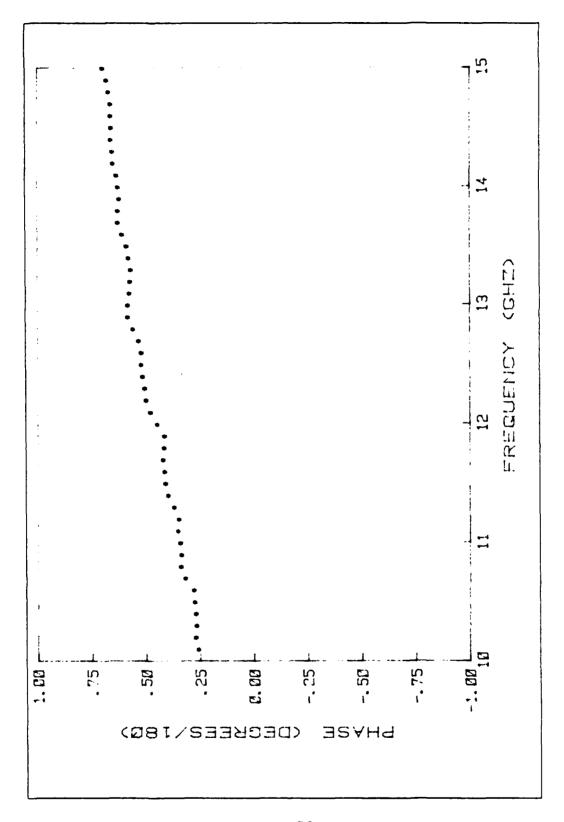


Figure 3.14 TARGET5 Phase Shift vs. Frequency

TABLE 7
TARGET5 Measured Data

| 10.10 00433 .24304 10.20 00431 .25591 10.30 00397 .25206 10.40 .00401 .25567 10.50 00436 .26109 10.60 .00452 .26315 10.70 .00410 .30425 | Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|---|--|--|--|
| 10.80 | 10000000000000000000000000000000000000 | 1771620671958762538555493341951346636276921632069444449955484573999499667778883777492163206944736236645554993341951346636276988857749216320698889777492163206988897774921632069888977749216320698889777492163206988897774921632069888997774921632069888997774921632069888997774921632069888997774921632069888997774921632069888999888999888998889988899888998889 | 167955435112:15532477.67248468249.1305717.9548136598585134374261348136598585143772452882371424848484981365717.954813652485524835524837724452892334247168344249813657344442359438868666666666666666666666666666666666 |

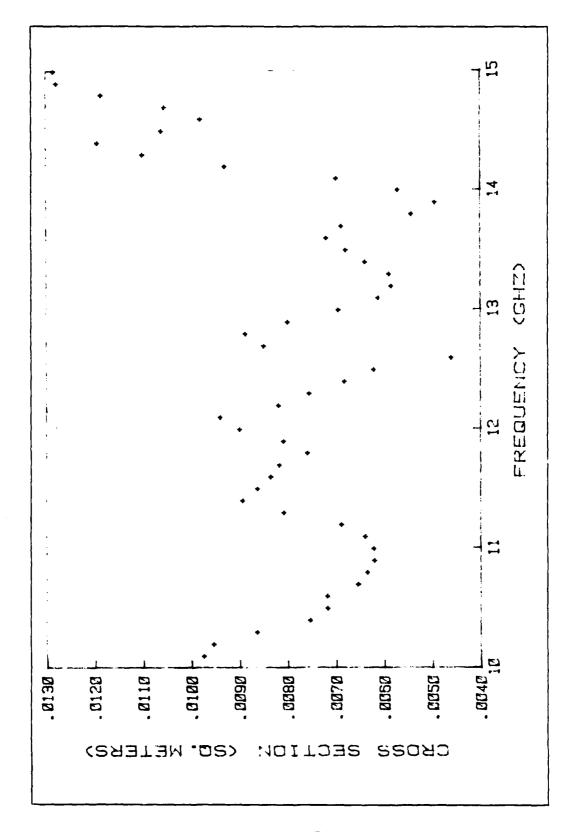


Figure 3.15 TARGET6 Cross-Section vs. Frequency

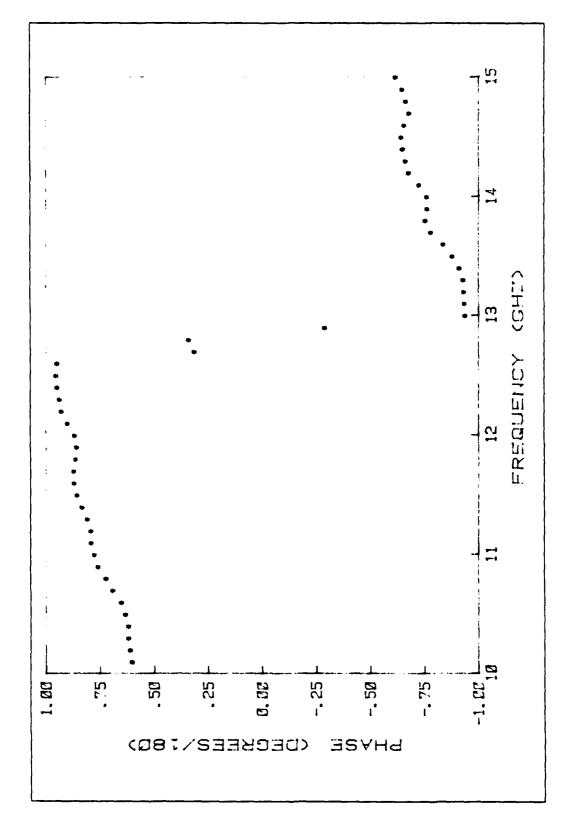


Figure 3.16 TARGET6 Phase Shift vs. Frequency

TABLE 8
TARGET6 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|---|--|---|
| 10 20 00 00 00 00 00 00 00 00 00 00 00 00 | 9947 9947 99947 9997 9997 9997 9997 999 | 6330183031544499997003305742378998057480526527622887301830315444999970033057423789980574805265276228853857777777788888888888899333057477777778976480762288574999988777777897648088888888888999333057777777897648067639228865739230574899887777789764866789999999887777789764866789999999999 |

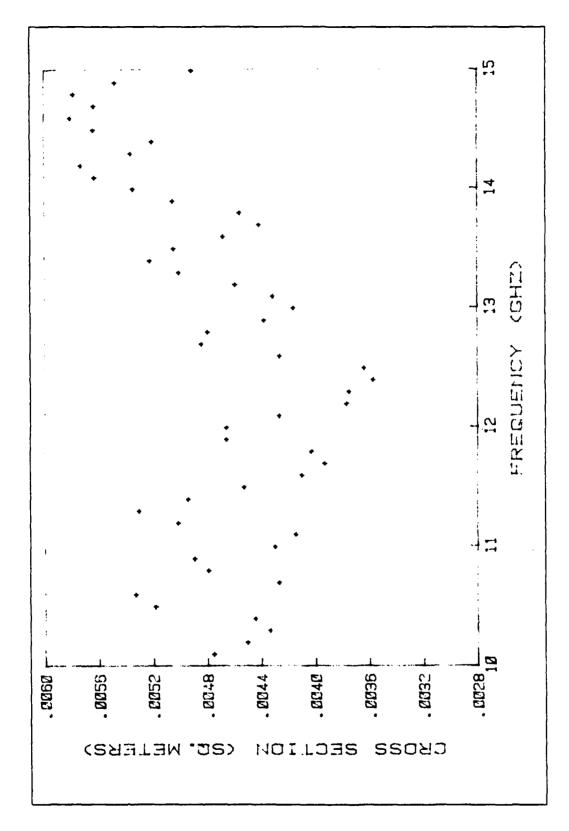


Figure 3.17 TARGET7 Cross-Section vs. Frequency

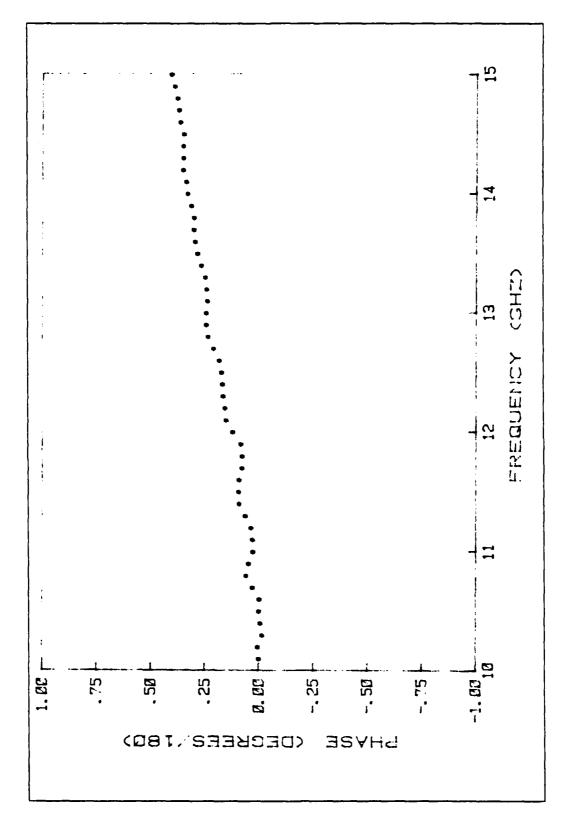


Figure 3.18 TARGET7 Phase Shift vs. Frequency

TABLE 9
TARGET7 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|------------------|---|---|
| 10 | .00473 .004474 .004471 .004477 .004477 .004477 .004473 .004453 .004473 | 72836524875135071113148505324977215106728845022652487513507111314159792485022653333333333333333333333333333333333 |

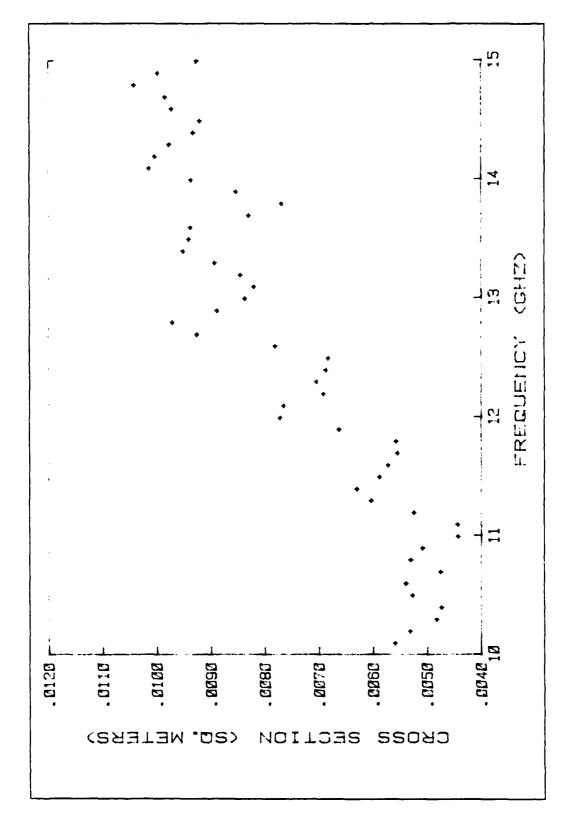
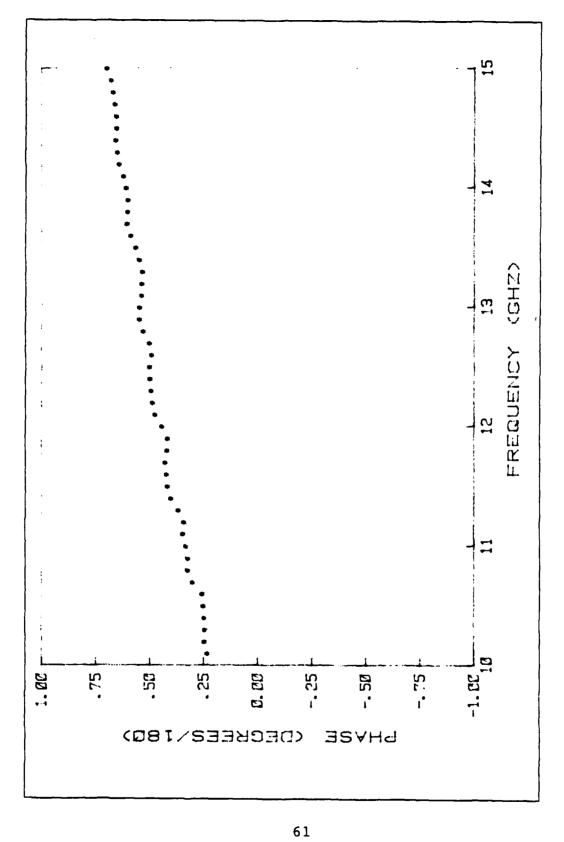


Figure 3.19 TARGET8 Cross-Section vs. Frequency



TARGET8 Phase Shift vs. Frequency Figure 3.20

TABLE 10
TARGET8 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|---|--|--|
| GH 1236000000000000000000000000000000000000 | sq. meters 5267814053778742692769691750522496541437087173686 00000000000000000000000000000000000 | Degree 182044917792152418841136722334429177921524142505941932334189449173523344649176475921524647764579527418846465533333333333333333333333333333333 |
| 14.90 15.00 | (00991 00920 | 66344 68325 |

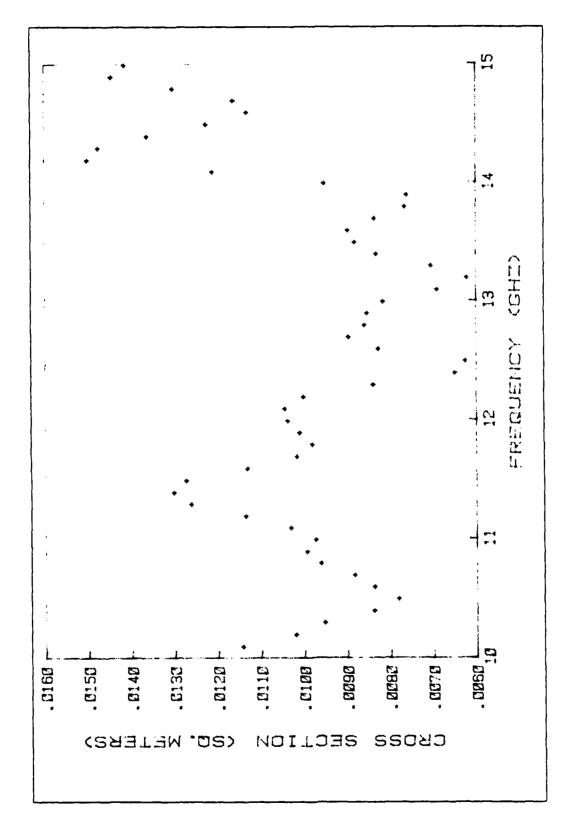


Figure 3.21 TARGET9 Cross-Section vs. Frequency

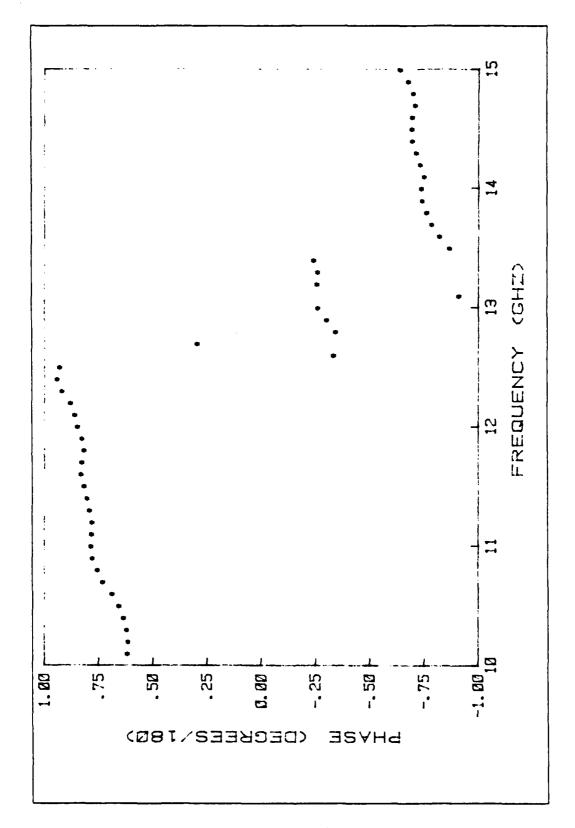


Figure 3.22 TARGET9 Phase Shift vs. Frequency

TABLE 11
TARGET9 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|------------------|--|---|
| 10 | 011011295724943090909090909090909090909090909090909 | 5265415840506142099028707564119100124777738753414315840584050888888889925641191 |
| 10 | 00875 009875 009877 009877 009877 009987 0010295 011255 011255 011297 011297 0110974 0109973 01099849 008899 | 715684 7159999999999797564 71597777777788149199999797564 715977777777778814572999914624477777777777888188888888889991462477777777777777777777777777777777777 |

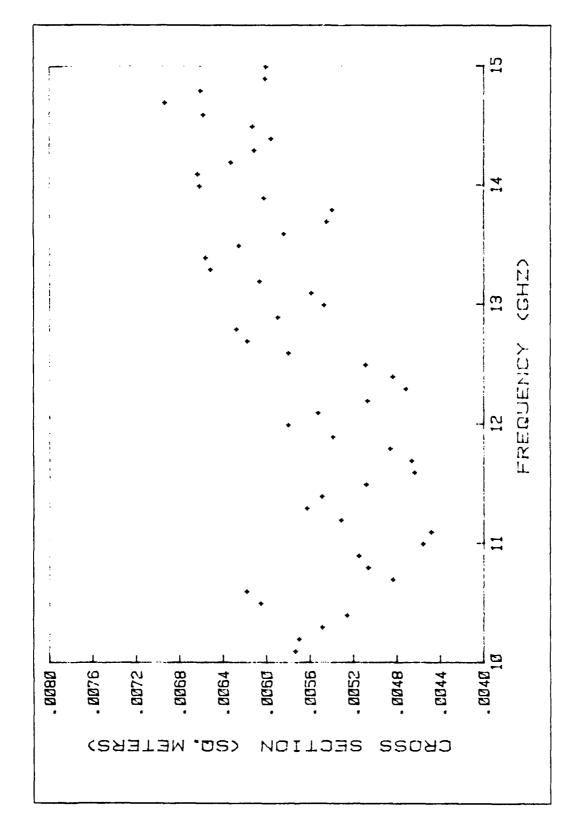


Figure 3.23 TARGET10 Cross-Section vs. Frequency

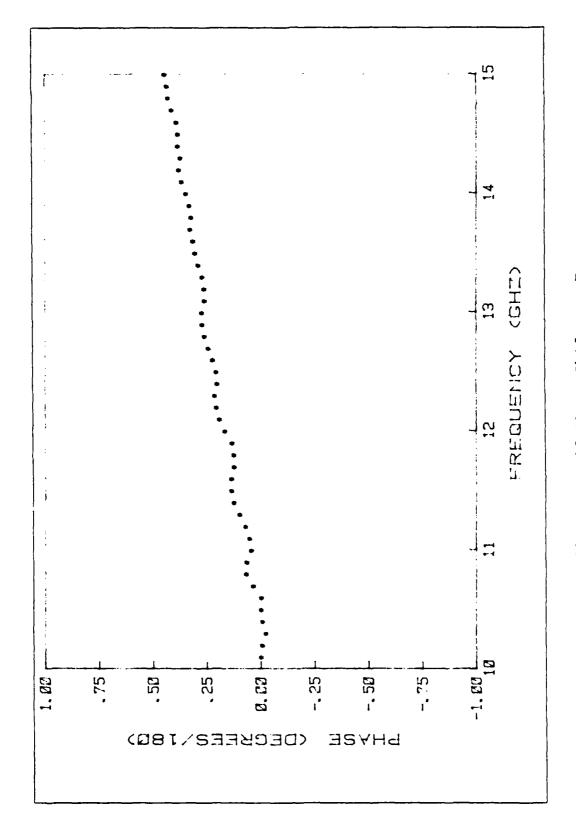


Figure 3.24 TARGET10 Phase Shift vs. Frequency

TABLE 12
TARGET10 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|---|--|---|
| 100.200 100.20 | 976325132358965133679481675574649331279999999999999999999999999999999999 | - 016413 - 02873 - |

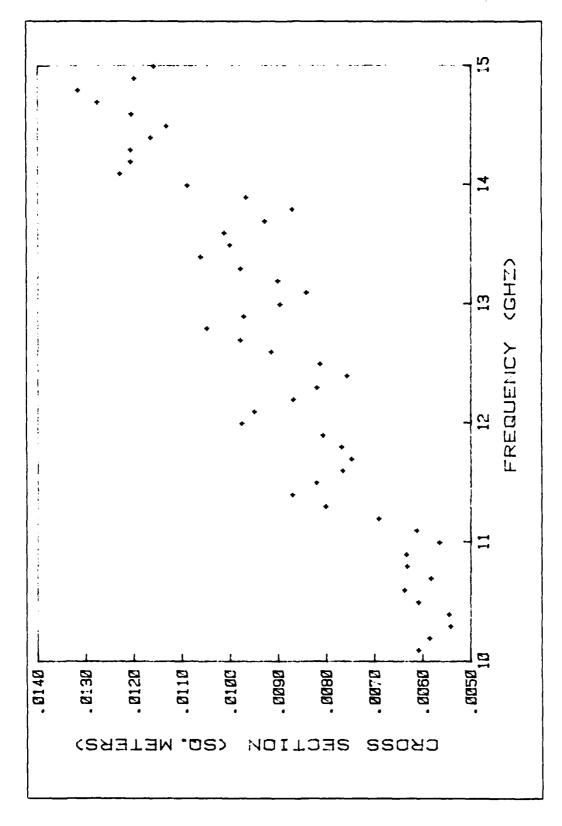


Figure 3.25 TARGET11 Cross-Section vs. Frequency

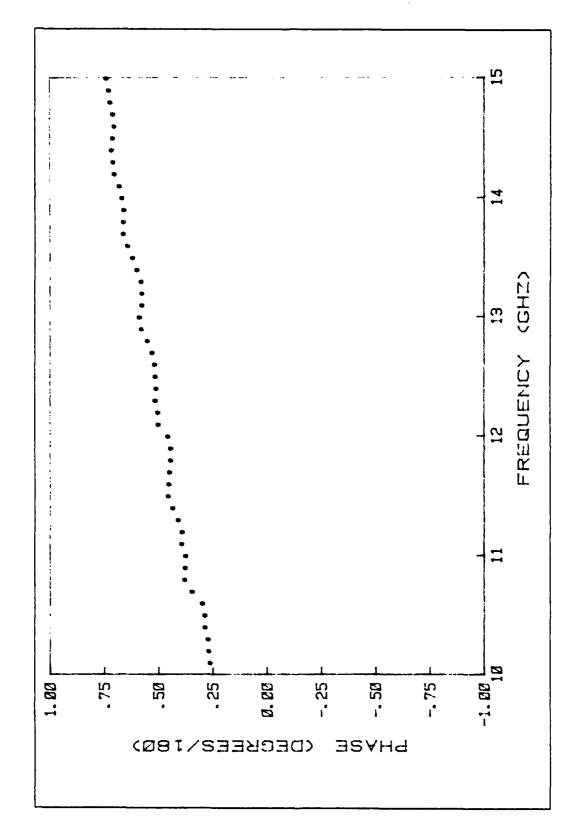


Figure 3.26 TARGET11 Phase Shift vs. Frequency

TABLE 13
TARGET11 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|-------------------------|--|--|
| 19 GH | sq. me 23 me 22 me | Degrees/180 2556944 2556944 2556944 2556944 2556944 2556944 2556944 2556944 2556944 2556944 2556944 2556944 2556944 2556944 25664 25664 25664 25664 25664 25664 25664 25664 25664 25664 25664 25664 25664 25664 25666 26 |
| 14.80 14.90 15.00 | .01310 .01193 .01151 | .70730 .71326 .72481 |

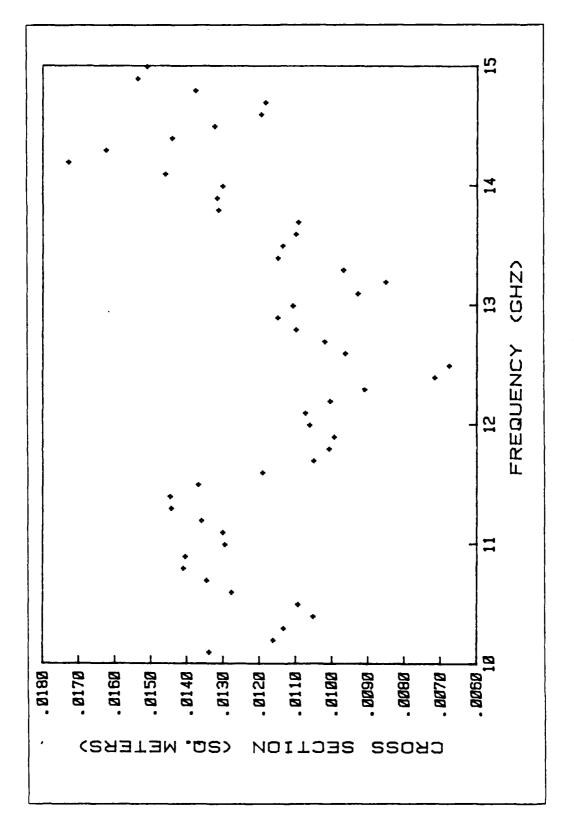


Figure 3.27 TARGET12 Cross-Section vs. Frequency

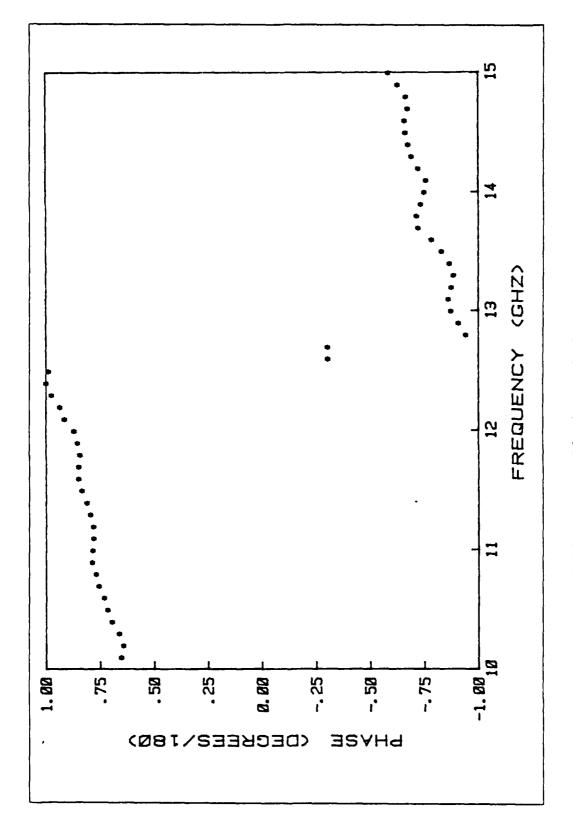


Figure 3.28 TARGET12 Phase Shift vs. Frequency

TABLE 14
TARGET12 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|---|--|--|
| 100 100 100 100 100 100 100 100 100 100 | 01323 0112424 01102637 01102637 01102637 01102637 01102637 01102637 01102637 0110263 | 03627753423517785607435510226407615889111475435017451777777778888889999833999878991114775435017459966666666666666666666666666666666666 |

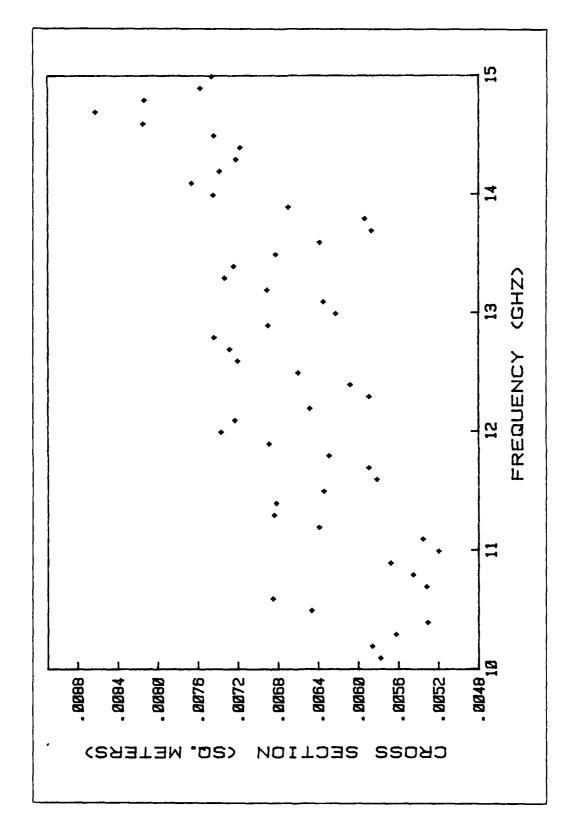
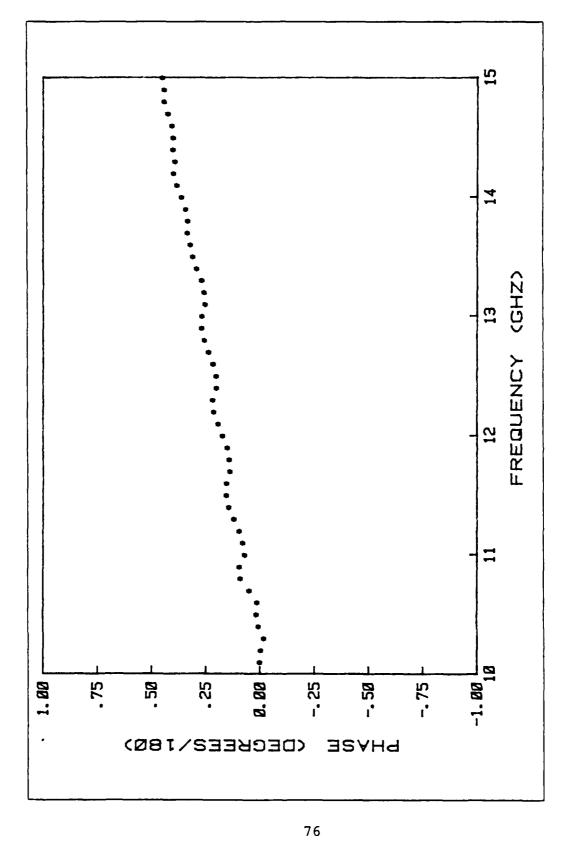


Figure 3.29 TARGET13 Cross-Section vs. Frequency



TARGET13 Phase Shift vs. Frequency Figure 3.30

TABLE 15
TARGET13 Measured Data

| Frequency GHz | Cross-Section sq. meters | |
|--|--|--|
| 10000000000000000000000000000000000000 | 0053391387 00053391387 00053391387 0006339 000639 | 7545915225222425425442544254425442542522222222 |

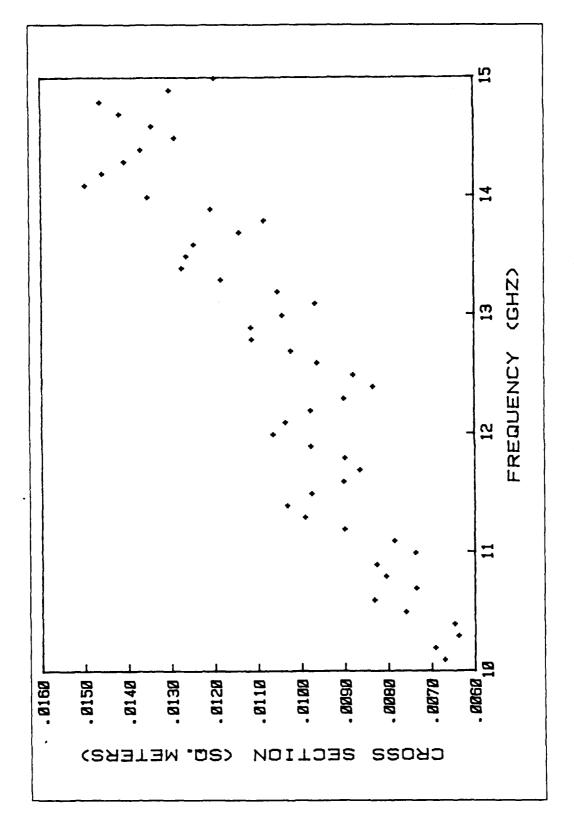


Figure 3.31 TARGET14 Cross-Section vs. Frequency

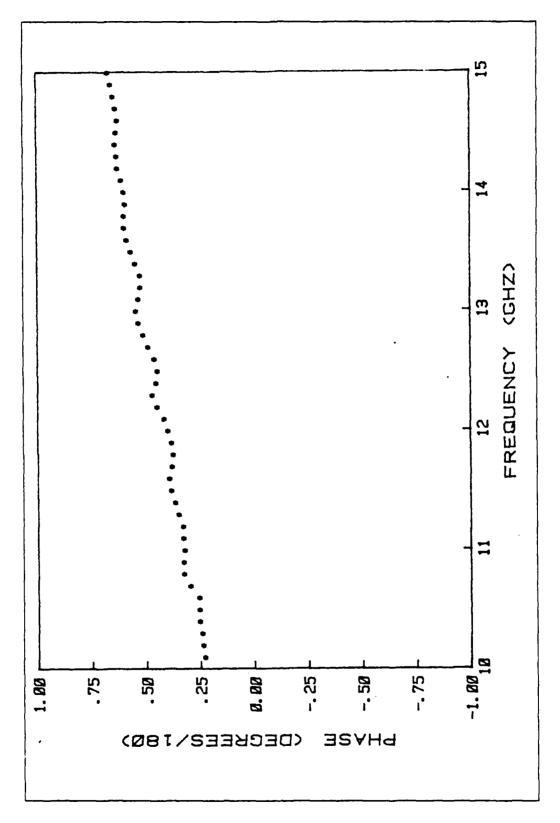
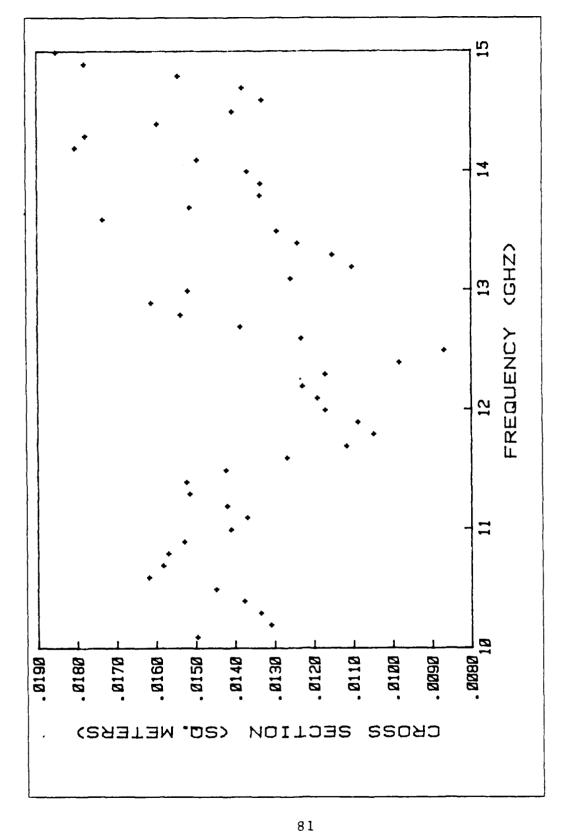


Figure 3.32 TARGET14 Phase Shift vs. Frequency

TABLE 16
TARGET14 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|--|--|--|
| 10000000000000000000000000000000000000 | ### 624 624 624 624 626 626 626 626 | 2224416946482166564217665944089590257142750440843133135582755555555555555556666666666666666666 |
| 15.00 | .01190 | .65430 |



TARGET15 Cross-Section vs. Frequency Figure 3.33

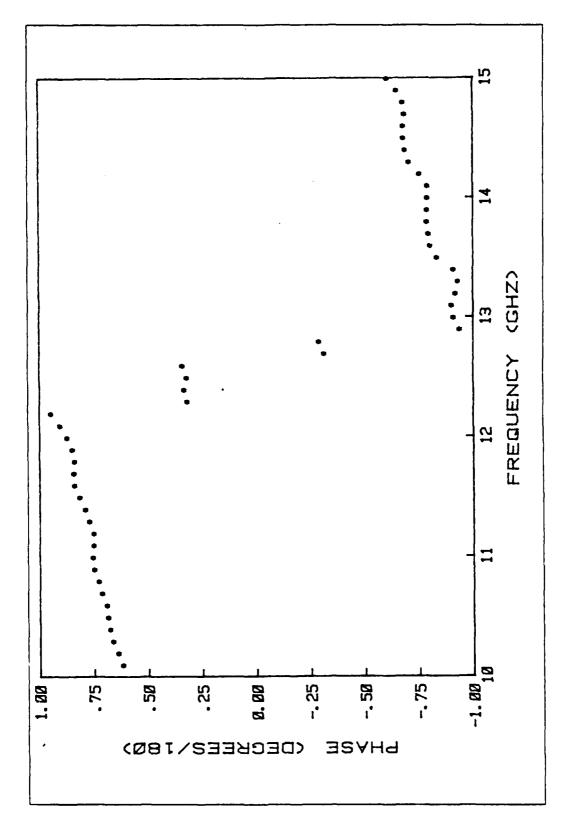


Figure 3.34 TARGET15 Phase Shift vs. Frequency

TABLE 17
TARGET15 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 | |
|---|---|--|--|
| 100.000.000.000.000.000.000.000.000.000 | 81688 91328 913364 913364 913364 913364 91455 91455 91455 9155 | 4992433953443881635946293497752443921730357891296455566189537777788888888933333351783035789129645555618518561856185618561856185618561856 | |

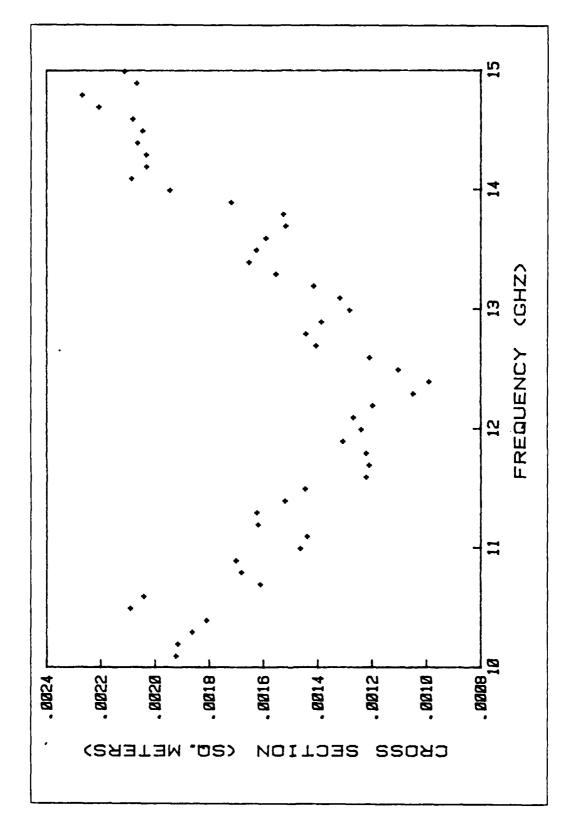


Figure 3.35 TARGET16 Cross-Section vs. Frequency

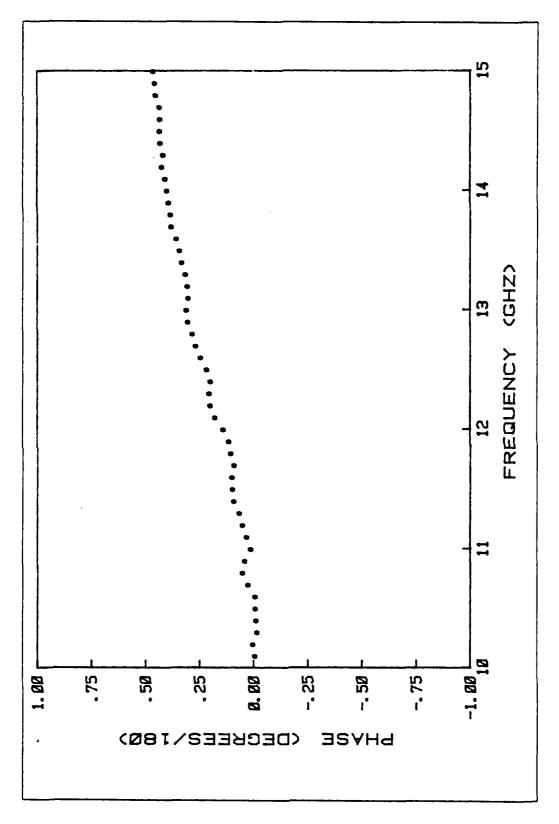


Figure 3.36 TARGET16 Phase Shift vs. Frequency

TABLE 18
TARGET16 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|--|---|----------------------|
| 10000000000000000000000000000000000000 | 9919898399795379699198993991989999999999999999999999999 | |

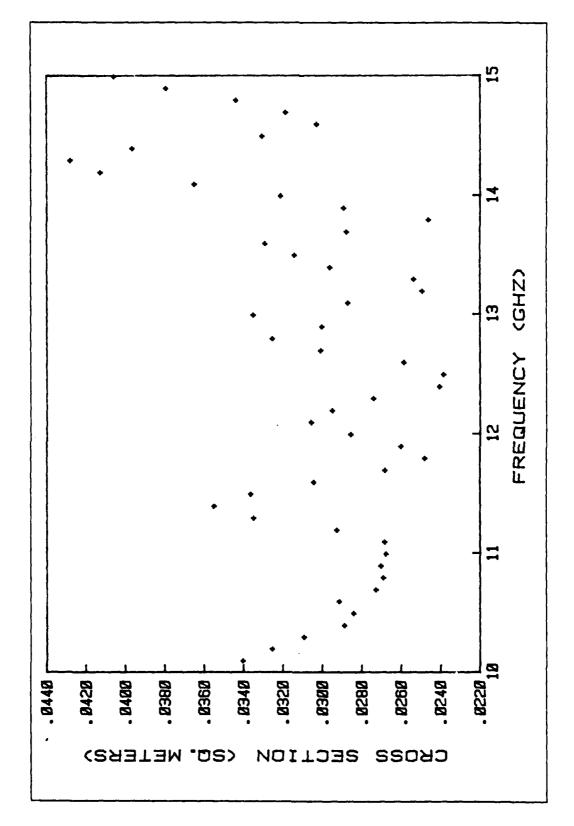


Figure 3.37 TARGET17 Cross-Section vs. Frequency

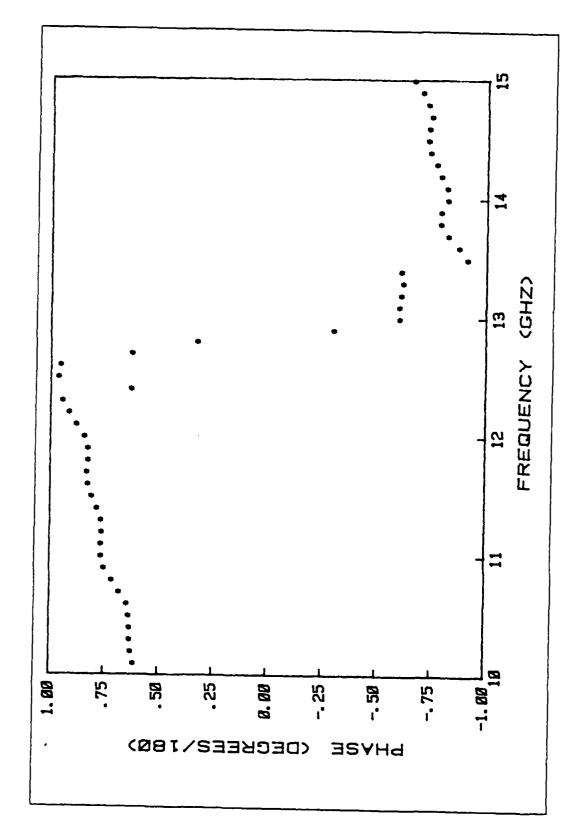


Figure 3.38 TARGET17 Phase Shift vs. Frequency

TABLE 19
TARGET17 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|--|---|--|
| 10000000000000000000000000000000000000 | 888726914817653487654923765579644316663133487654923357796914817633334876654923357796914817666513348766549233577964431666313348765349293357796443166631333487665492335779644316663133333333333333333333333333333 | 704666989682968296475995844296964453889594467538662974286964538896984767777788888888969447297777888388888969447579427777777778888888889694747777777777777777 |
| 14.90 15.00 | . 04940 | 67494 |

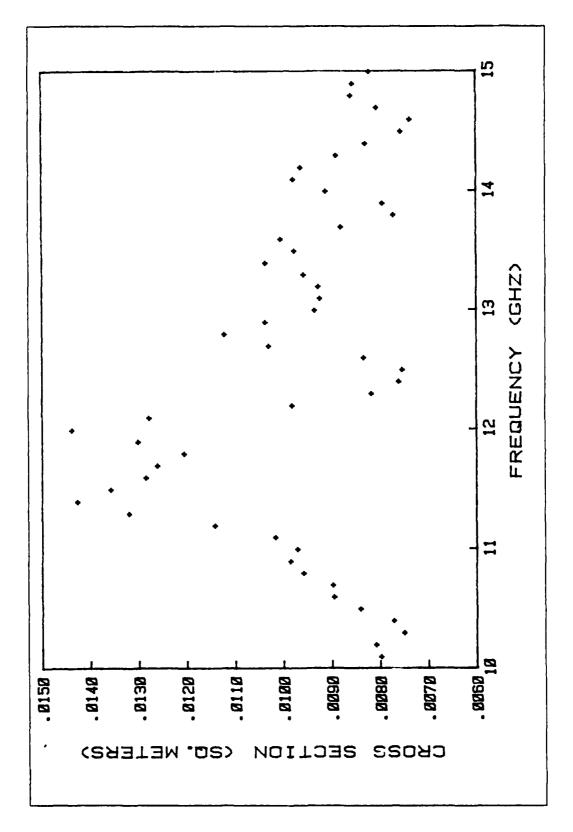


Figure 3.39 TARGET18 Cross-Section vs. Frequency

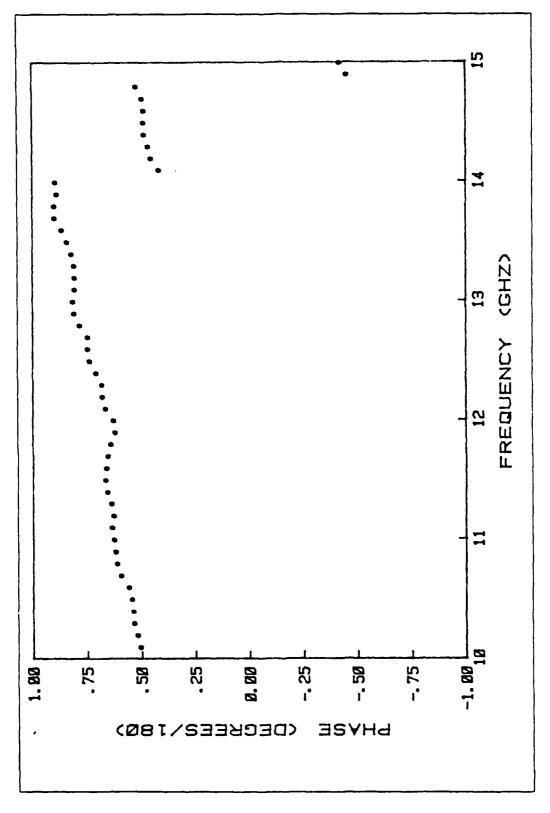


Figure 3.40 TARGET18 Phase Shift vs. Frequency

TABLE 20
TARGET18 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|--|--|--|
| 10.230 10 | 0007339 00074559 00074559 00074559 00077339 00077339 00077339 00097747 0013129 | ###################################### |

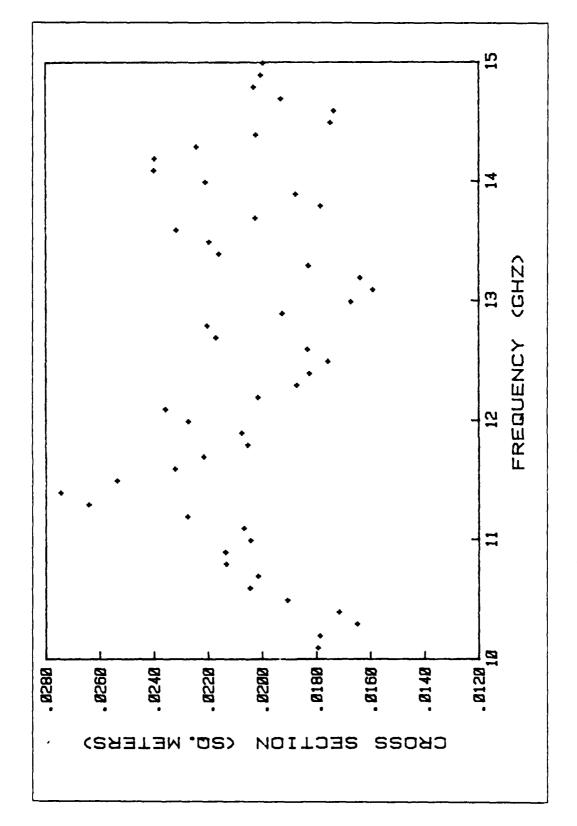


Figure 3.41 TARGET19 Cross-Section vs. Frequency

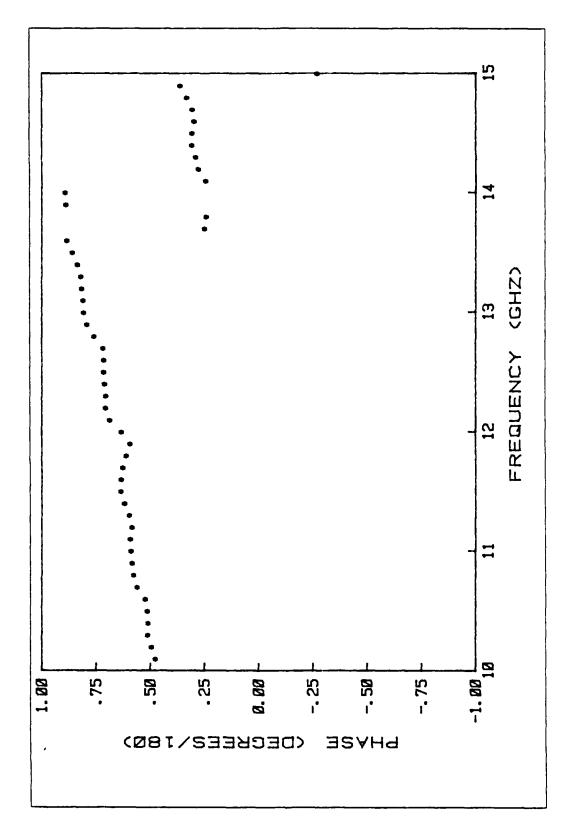
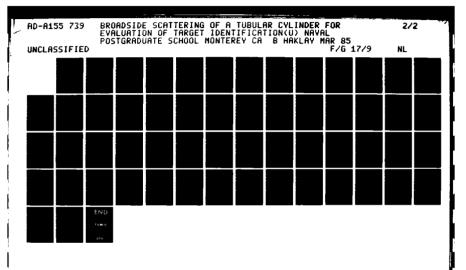
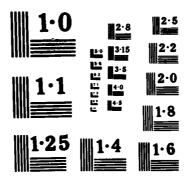


Figure 3.42 TARGET19 Phase Shift vs. Frequency

TABLE 21
TARGET19 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|--|---|---|
| 10000000000000000000000000000000000000 | 348454315264924052315495493408678754247740163993408226752163941639944086787542477401639934226752200664922405233993408678213016399349934993499349934993499349934993499 | 5917993311113434745518154693343693336947761489810561863944491111347455181546933443833333694776148981056678644996566867786139978135666667777778888882288222222223356 |
| | . was a series of | |





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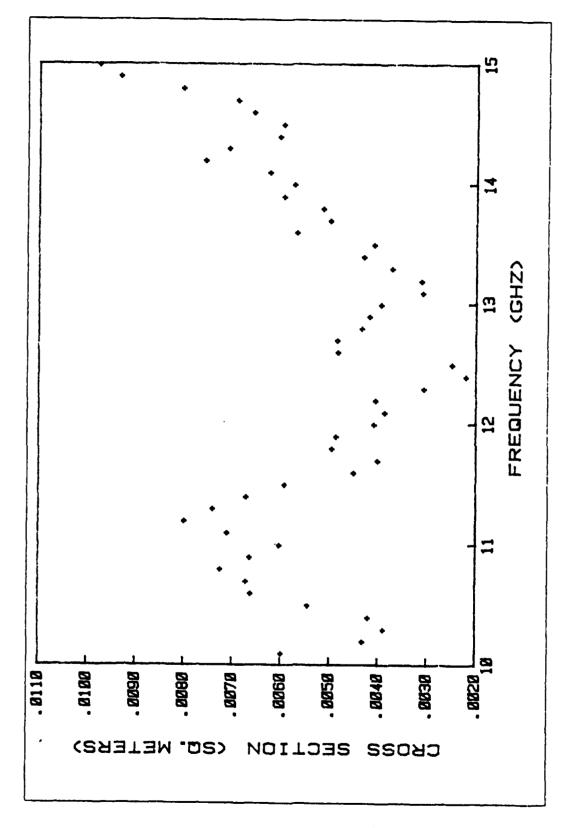


Figure 3.43 TARGET20 Cross-Section vs. Frequency

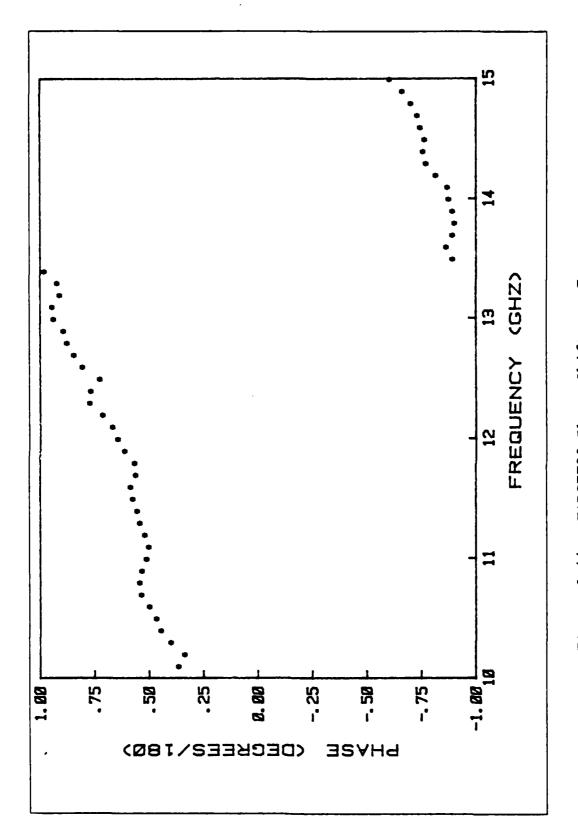


Figure 3.44 TARGET20 Phase Shift vs. Frequency

TABLE 22
TARGET20 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|--|---|---|
| 10000000000000000000000000000000000000 | 255238555777424564598289953278829366421388878127924759 5438135557777424564598828995327728829366421388878127924759 888888876577776884988999147722188888 88888888888888888888888888888 | 9888819191949649270562597360191495507005448397472520795868191949649270562597360191495507005448397472520733344828755564498259777778887929987877777778887999888777777786672299060770111111111111111 |

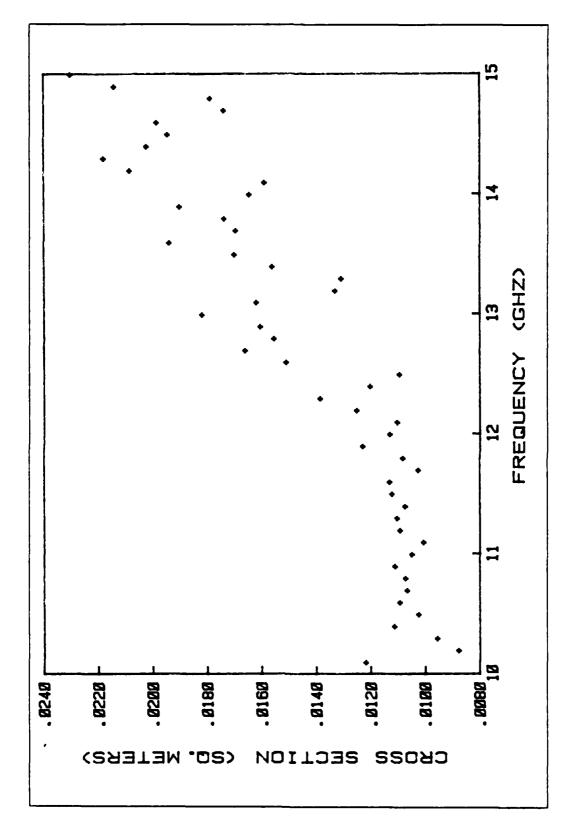
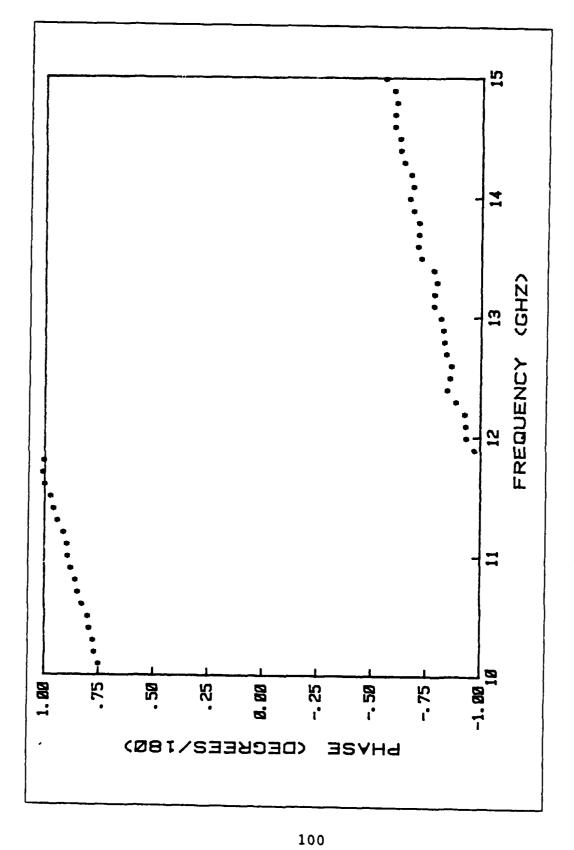


Figure 3.45 TARGET21 Cross-Section vs. Frequency



TARGET21 Phase Shift vs. Frequency Figure 3.46

TABLE 23
TARGET21 Measured Data

| Frequency GHz | Cross-Section sq. meters | Phase Degrees/180 |
|--|--|--|
| 10000000000000000000000000000000000000 | 0126644 0126644 0126644 012669 011001001001001001001001001001001001001 | 187537 1875387 1875 |

IV. DATA ANALYSIS

A. ANALYSIS OF EXPERIMENTAL DATA

The experimental data on the back scattered cross section of a tubular cylinder obtained in the Scattering Laboratory as described in Chapter III are analyzed. The cross section under the circumstances described in Chapter III is a function of three parameters:

- (1) The length of the cylinder (2h)
- (2) The diameter of the cylinder (2a).
- (3) The frequency of the incident wave (f). This is shown in equation 4.1

$$\sigma = G(a,h,f) \tag{4.1}$$

The experimental data as shown in Figures 3.5 to 3.46 shows the dependence of the cross section on frequency for cylinders of fixed lengths and diameters.

To isolate the dependence of the back scattering cross section on the length of a cylinder, the constant parameteres should be the cylinder diameter and the frequency. This was achieved by using five cylinders with the same diameter and taking the cross section of each one of them at the same frequency. Table 24 shows the targets used for each diameter.

The frequencies checked were 10.1,11,12,13,14 and 15 GHz and the results are given at the end of this chapter in Figures $4.1,\ 4.2,\ 4.3$.

In Figure 4.1 for 2a=0.375", one can see that the curve has tilt at f=10.1 when 2h=2.7; for f=11 the tilt occures at 2h=2.5 and for f=12 this happens at 2h=2.25. The line seems

TABLE 24
Targets with Constant 2a

| 2h 2a | 0.375" | 0.5" | 0.75" |
|-------|----------|----------|----------|
| 2.0 | TARGET1 | TARGET2 | TARGET3 |
| 2.25 | TARGET4 | TARGET5 | TARGET6 |
| 2.5 | TARGET7 | TARGET8 | TARGET9 |
| 2.75 | TARGET10 | TARGET11 | TARGET12 |
| 3.0 | TARGET13 | TARGET14 | TARGET15 |

to have the same slope in all these graphs near this point. The same phenomenon appears for 2h=2.75 at f=11; 2h=2.5 at f=12 and 2h=2.25 at f=13. In Figure 4.2 for 2a=0.5" the same phenomenon occurs at 2h=2.75 and f=12, 2h=2.5 and f=13 and 2h=2.25 and f=14. For 2a=0.75" in Figure 4.3 this occurs at 2h=2.75, f=11 and 2h=2.4, f=12; and another at 2h=2.5, f=13 and 2h=2.25, f=14.

This phenomenon leads to the assumption that the back scattering cross section of a tubular cylinder has a dependence not on h by itself but on combination of f and h. For those points mentioned above, the product 2hf is shown in Table 25.

Thus the cross section of a tubular cylinder can be written as a function of hf or kh where $k=2\pi f/c$ and is shown in equation 4.2 which is identical to equation 4.1 .

$$\sigma = G_1(a, f, hf) = G_2(a, k, kh)$$
 (4.2)

B. COMPARISON BETWEEN MEASUREMENTS AND THEORY

From equations 2.30 one can see that σ , when properly normalized, (e.g. divided by ah or multiplied by k^2) depends only on two variables which are combinations of a, h and k. This dependence is shown in equation 4.3, where l_1 =kh and l_2 =ka. This equation can also be written as 4.4.

TABLE 25
2hf for Discontinuity Points

| 2a | f | 2h | 2hf |
|-------|------|------|-------|
| | | | |
| • | 10.1 | 2.7 | 27.25 |
| 0.375 | 11 | 2.5 | 27.5 |
| | 12 | 2.25 | 27.0 |
| | | | |
| | 11 | 2.75 | 30.25 |
| 0.375 | 12 | 2.5 | 30.0 |
| | 13 | 2.25 | 29.25 |
| | | | |
| | 12 | 2.75 | 30.25 |
| 0.5 | 12 | 2.75 | 30.0 |
| | 13 | 2.25 | 29.25 |
| | | | |
| 0.75 | 11 | 2.75 | 30.25 |
| | 12 | 2.4 | 28.8 |
| | | | |
| 0.75 | 13 | 2.5 | 52.5 |
| | 14 | 2.25 | 31.5 |
| | | | |

$$\sigma/ah=F(1_1,1_2) \tag{4.3}$$

$$\sigma/ah=F_1(l_1,l_2/l_1)=F_1(ka,h/a)$$
 (4.4)

For cylinders with a constant h/a, σ/ah can be plotted as a function of ka on the same graph. This effectively expands the frequency range over which data can be obtained using a cylinder. Table 26 shows the targets used for this purpose.

TABLE 26 Cylinders with the Same h/a

| Target name | Length | Diameter | h/a |
|-------------|--------|----------|-----|
| TARGET16 | 1.5" | 0.375" | 4 |
| TARGET2 | 2" | 0.5" | 4 |
| TARGET18 | 2.5" | 0.625" | 4 |
| TARGET15 | 3" | 0.75" | 4 |
| | | | |
| TARGET4 | 2.25" | 0.375" | 6 |
| TARGET14 | 3" | 0.5" | 6 |
| TARGET19 | 3.75" | 0.625" | 6 |
| TARGET17 | 4.5" | 0.75" | 6 |
| | | | |

Figure 4.4 is the graph for h/a=4 and Figure 4.5 for h/a=6. The ka range covered by each target is shown on Table 27.

TABLE 27 ka Range Covered by Each Target

| Target Name | h/a | minimum ka | maximum ka |
|-------------|-----|------------|------------|
| TARGET16 | 4 | 1.01 | 1.50 |
| TARGET2 | 4 | 1.34 | 2.00 |
| TARGET18 | 4 | 1.68 | 2.49 |
| TARGET15 | 4 | 2.01 | 2.99 |
| | | | |
| TARGET4 | 6 | 1.01 | 2.00 |
| TARGET14 | 6 | 1.34 | 2.00 |
| TARGET19 | 6 | 1.68 | 2.49 |
| TARGET17 | 6 | 2.01 | 2.99 |

The overlapping points as shown on these graphs show that the experimental data are in agreement with the overall

shape. However, the small variations in the overlapping regions do not seem to fall at the same places. The reason for this discrepancy were discussed as measurement errors in Chapter III.

The overall shape is used to compare with theoretical predictions obtained by Professor Lee at the Naval Postgraduate School through solving equations 2.30 and 2.31 [Ref. 13]. This data is shown in Figure 4.6 for h/a=4 and in Figure 4.7 for h/a=6. Theoretical values and experimental data are plotted together in Figures 4.8 and 4.9. In this figures one can see that there is very good agreement up to the point where ka is about 1.9. From that point to ka equal approximately 2.5 the minima and maxima of the experimental curve is shifted in ka by about 0.08.

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The reason for the shifting and the disagreement these points can be explained by the fact that there is wall thickness in the measured targets while the theory assumes infinitesimal thickness. The points of ka=1.9 and ka=2.5 are special because they correspond to the first two cutoff frequencies for the cylindrical waveguide modes: ka=1.8415 is the cutoff point for H_{11} mode and ka=2.4046 is the cutoff point for E_{01} mode [Ref. 14]. At ka=1.8415 the wave starts to propagate without attenuation inside the cylinder. Since this value is determined by the inner diameter of the cylinder, the wall thickness cause the phenomenon to occur at a higher frequency and so produces the shift in the minima and maxima. Above ka=2.4, E_{01} mode adds to the H_{11} mode and the total field inside the cylinder is the vectoral sum of those two modes. These effects were included in the theoretical calculation but the wall thickness was not.

The measured phase shift is shown in Figures 4.10 and 4.11 and the theoretical phase shift is given in Figures 4.12 and 4.13. The comparison as shown in Figures 4.14 and 4.15 shows agreement between the theory and the experiment.

The constant phase shift in the overall curve is due the calibration of the system. It was done with a 3.187" sphere while the targets are 0.375" to 0.75" in diameter. That caused small phase shift because the targets were not at the same height as the calibration sphere. At ka=2.75 where the phase shift is near ± 180 degrees, the average of ± 180 degrees produced the data points near 180(m-2n)/m degrees which appear in Figures 4.14 and 4.15 . Here $0 \le n \le m$ and m is the number of averages taken.

The comparison between the theory and the measured data leads to conclusions that can be used for future work in this project, that will discused in Chapter V.

C. RESULTS FOR CYLINDERS WITH FINS

As the first step for future work on employing target identification scheme through the cross section as a function of the frequency, two cylinders with fins attached have been measured (TARGET20 and TARGET21). The results as shown in Figure 3.43 and 3.44 are compared with the cylinder of the same length and diameter, TARGET15. The fins are like reflectors in TARGET20 and like corner reflectors in TARGET21. The comparison shows that the back scattered cross section of TARGET20 is much smaller then TARGET15. reason might be that the axial current that flow on the fins cause field that add vectorally to the field created by the current flowing on the cylinder and in this spatial case it happens to be added destructively. A detailed study on the surface current distribution on the cylinder without fins may lead to explanations about how the fins change surface current flow on the cylinder.

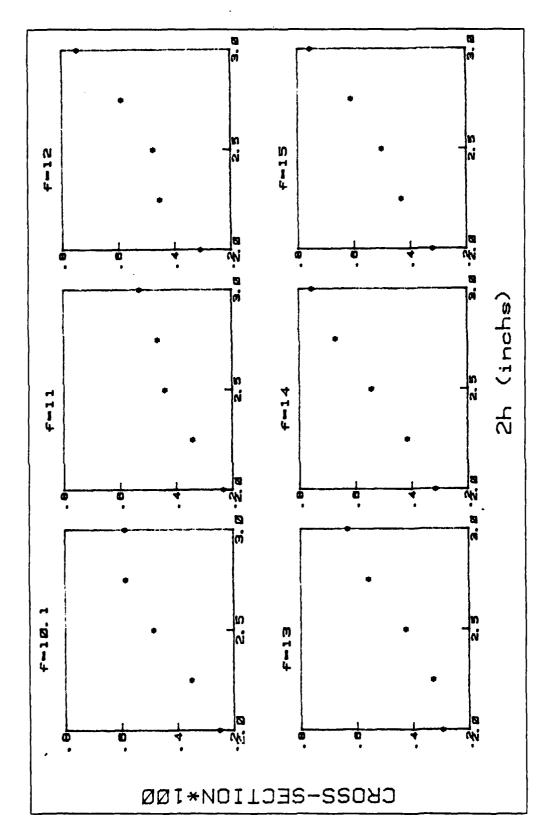


Figure 4.1 Length Dependence of Cross Section for 2a=0.375.

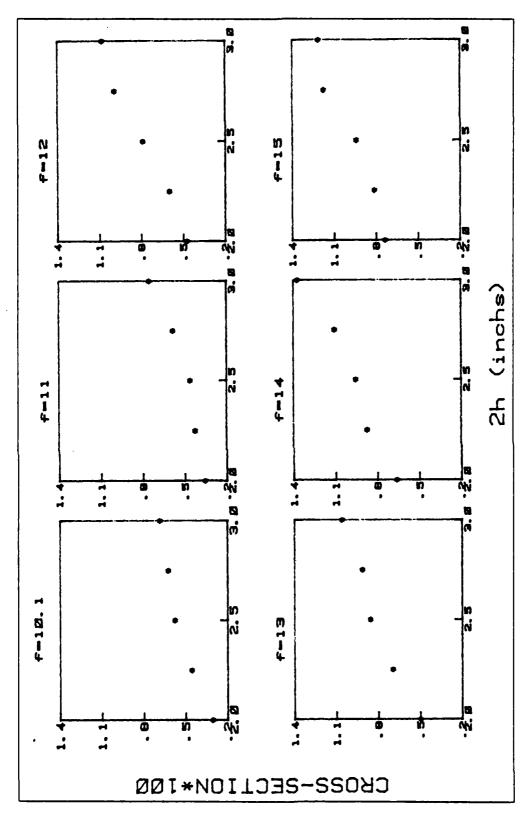
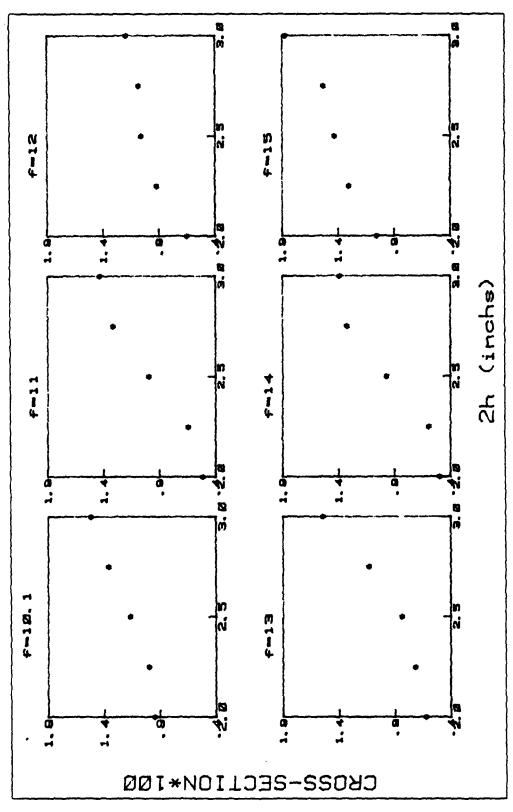
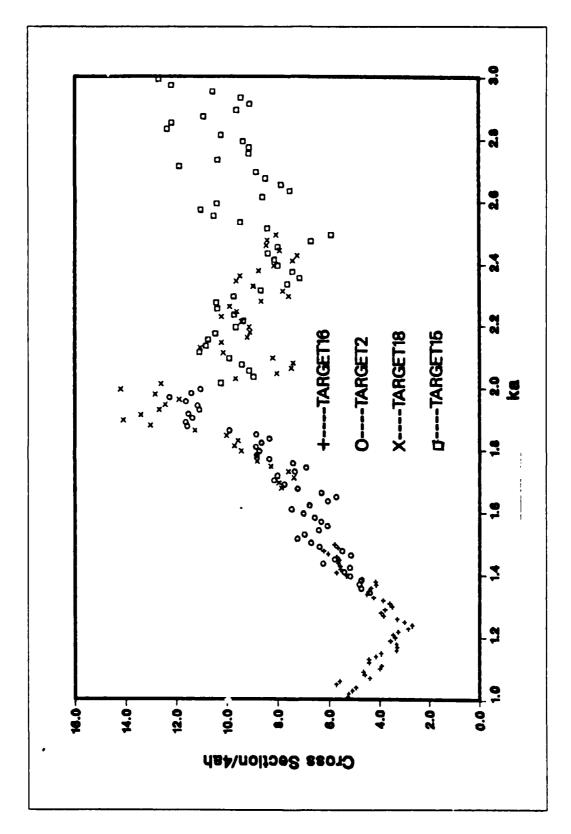


Figure 4.2 Length Dependence of Cross Section for 2a=0.5.

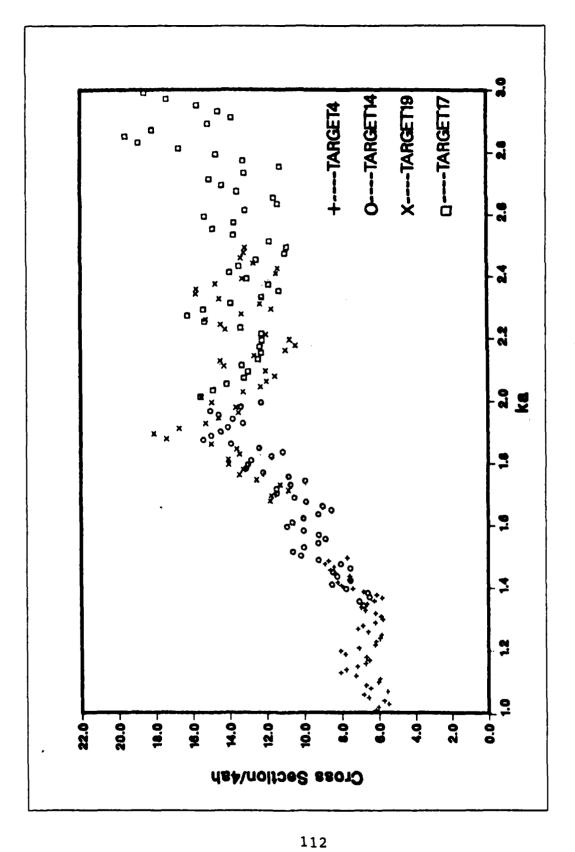


を全量が必要がある。 国立ののので 電影を行

Figure 4.3 Length Dependence of Cross Section for 2a=0.75.



Measured Cross Section/4ah vs. ka for h/a=4 Figure 4.4



Measured Cross Section/4ah vs. ka for h/a=6 Figure 4.5

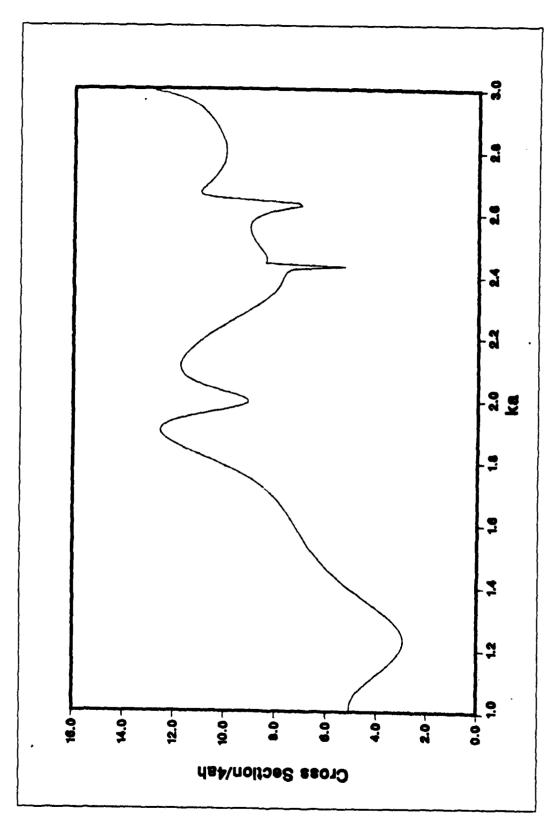


Figure 4.6 Theoretical Cross Section for h/a=4

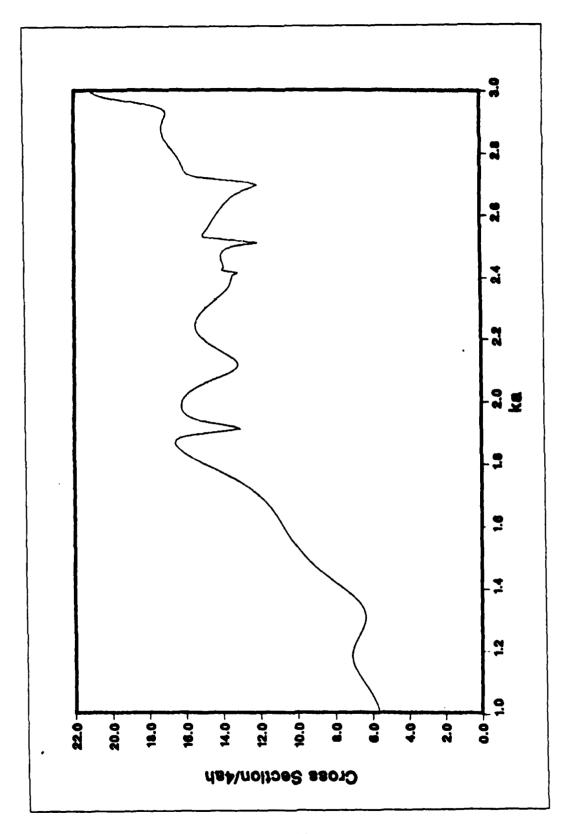


Figure 4.7 Theoretical Cross Section for h/a=6

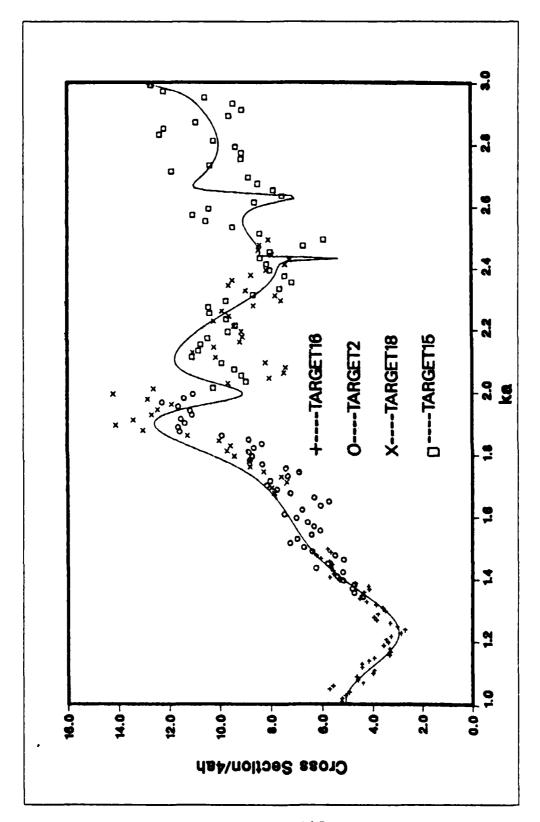


Figure 4.8 Comparison of Cross Section Between Theoretical & Experimental Data for h/a=4.

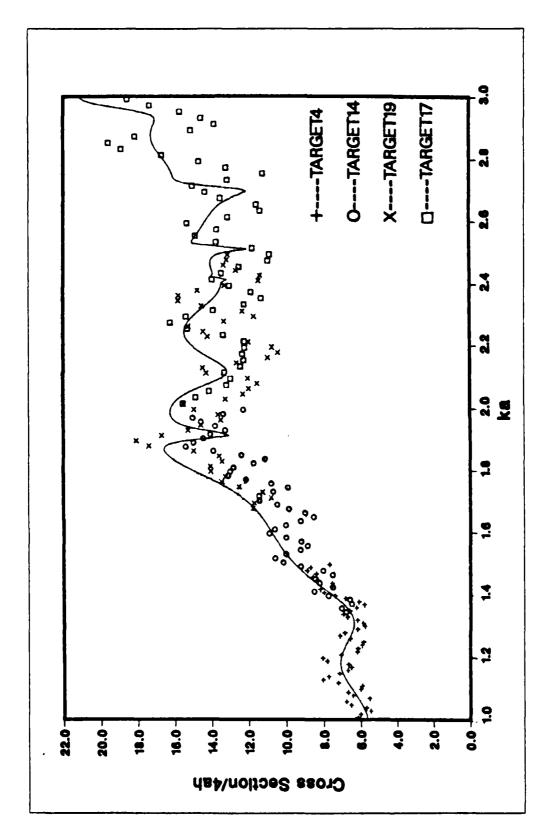


Figure 4.9 Comparison of Cross Section Between Theoretical & Experimental Data for h/a=6.

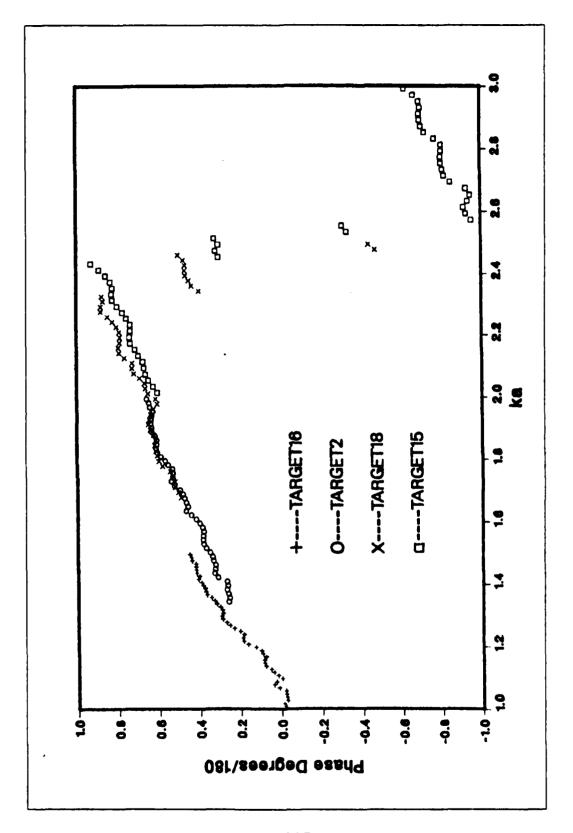


Figure 4.10 Measured Phase Shift vs. ka for h/a=4

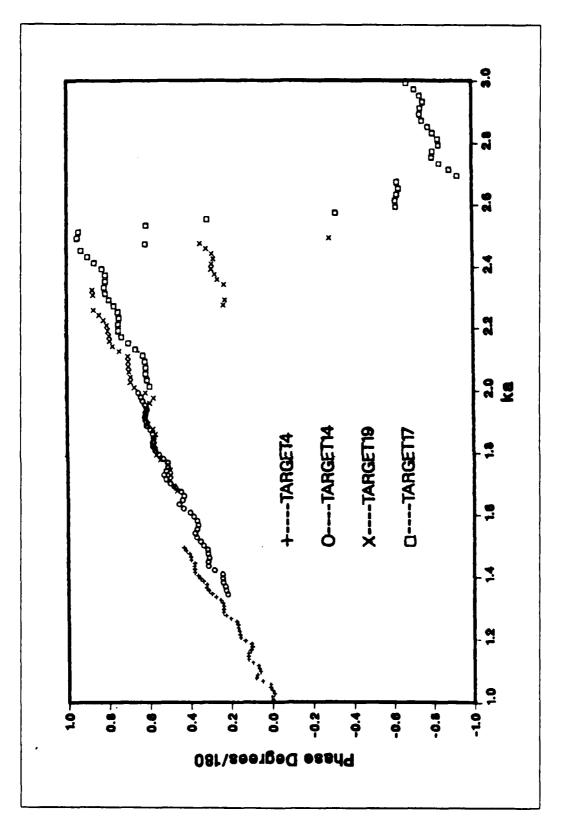


Figure 4.11 Measured Phase Shift vs. ka for h/a=6

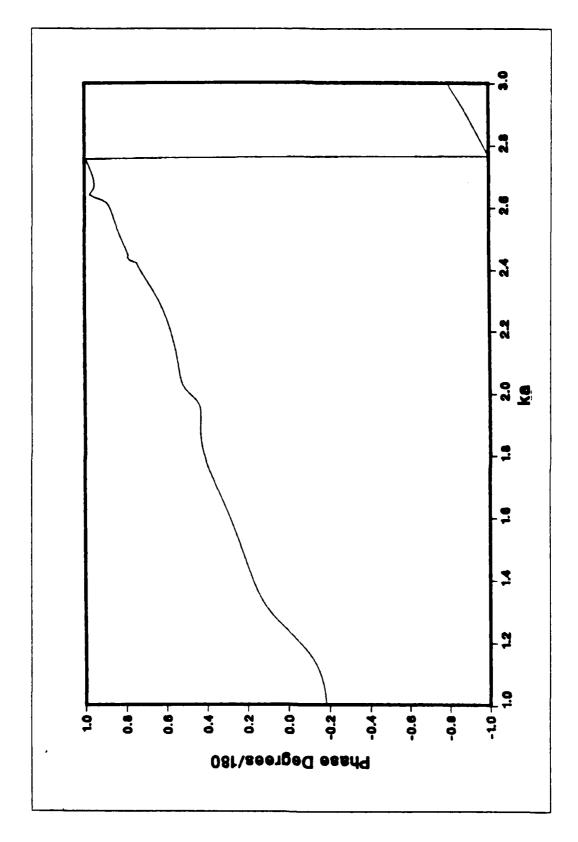


Figure 4.12 Theoretical Phase Shift for h/a=4

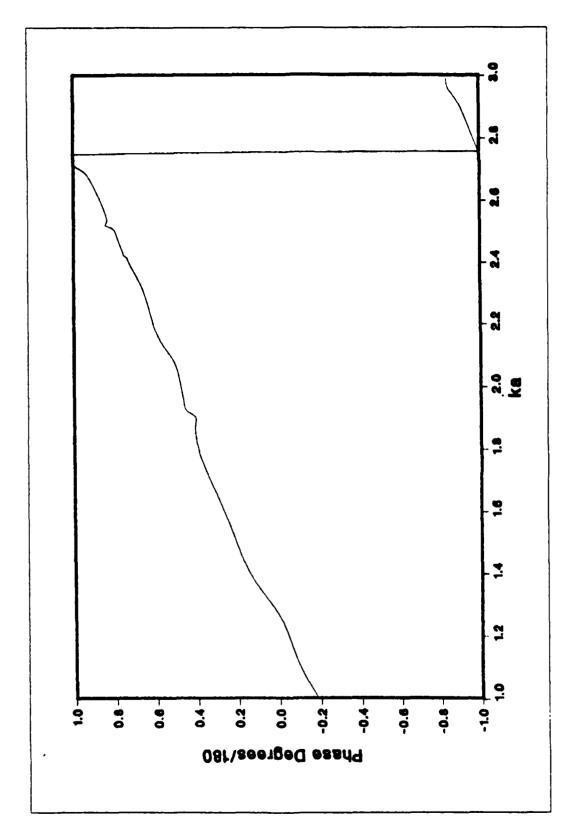


Figure 4.13 Theoretical Phase Shift for h/a=6

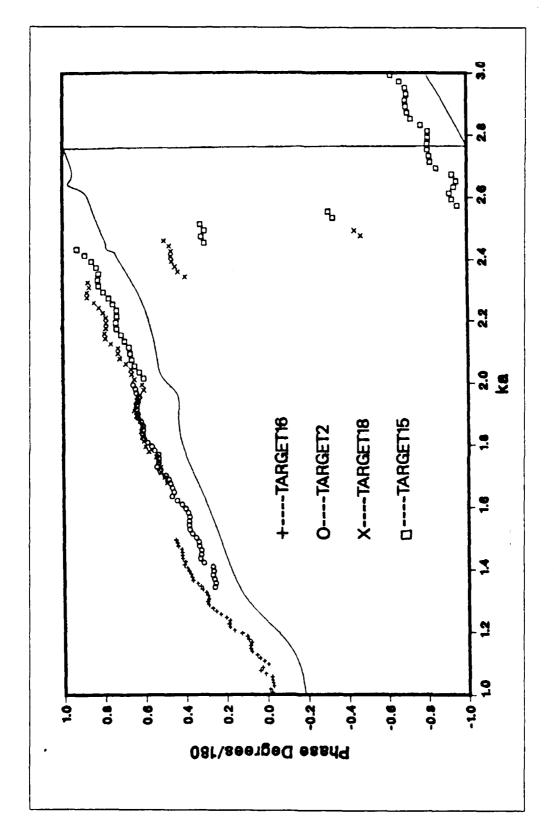


Figure 4.14 Comparison of Phase Shift Between Theoretical & Experimental Data for h/a=4.

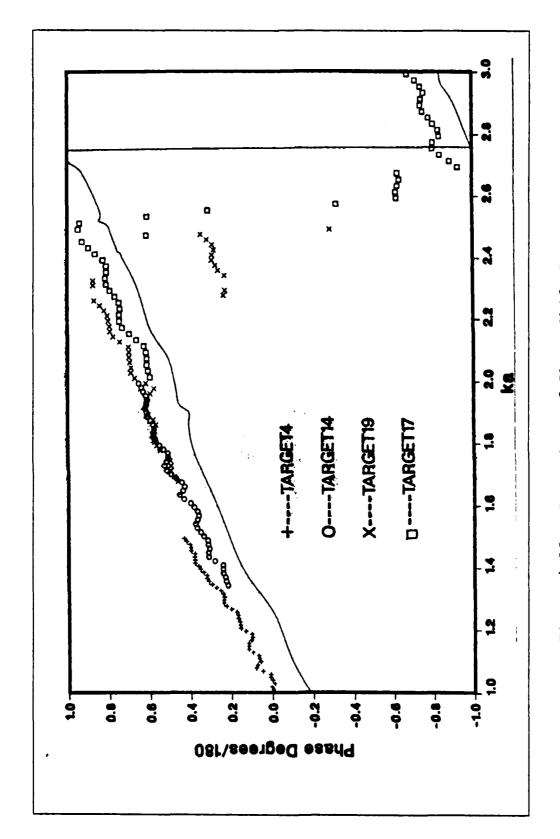


Figure 4.15 Comparison of Phase Shift Between Theoretical & Experimental Data for h/a=6.

V. SUMMARY

The exact solution to the scattering by a tubular cylinder from the broadside was presented. From that solution one can see that at this aspect angle the scattered field dependes only on the axial surface current flowing in the z direction. The exact solution was used to develop theoretical predictions to the broadside back scattering cross section and phase shift of a finite tubular cylinder. Measurements of cross section and phase shift have been done and the results were analyzed and compared to the theoret-The comparison shows that the results of the ical data. measurement and the theory are in very good agreement. Both theoretical and experimental data shows that the frequency range that has been chosen is good because the back scattering cross section is about fifteen times bigger then the cylinder dimensions, or equivalently, the cross section in the optical region. The increment in the cross section made the signal easier to be detected and analyzed.

These results as well as the problems that have been arised will be used for further work on this project in developing target identification schemes through the observation of the back scattered field from a target.

A. KNOWN PROBLEM AREAS

When more complicated bodies than the tubular cylinder are used as targets, the only way to study thier back scattered field is by using the measured data. This is why the most important problem that has to be addressed is to understand the discrepancies between theory and experiment that have been discovered. As a first step, the effects of the support should be studied by using different supports.

Another problem that discovered is the error in the average procedure for phase shift near ±180 degrees. That can be solve with a proper averaging procedure. Because of the system noise in the lower frequency range the measurements were done over the range of 10 to 15 Ghz and required the use of several scaled cylinders to expand the frequency range. This introduced errors due to the differences in the inner to outer diameters ratio. By achieving stability in the frequency response of the receiver, one cylinder can be used for a larger frequency range and thus all the parameters will remain constant for all the frequencies.

B. WHERE FUTURE WORK IS NEEDED

To understand the effects of cylindrical waveguide modes on the back scattering of a tubular cylinder, the solid cylinder should be studied. The closed ends of the cylinder introduce current in the transverse directions but prevent the wave from propagating inside the cylinder. Adding fins to the cylinder to investigate their effects on the induced surface current on the cylinder should be carried out next. The effects of the dimensions of the fins and their position on the cylinder, on the scattered fields are the most interesting.

Tests on models of real targets are the last step in the evaluation of this scheme and different aspect angles should be studied. For real-life targets, the wavelength should be scaled as the ratio between the models and the targets. That means production of radars that will work at frequencies at a few hundreds of MHz. This system should include a computer that has the ability to store data about the parameters of different targets of interest and by comparison between measured data and stored data the target will be identified.

APPENDIX A ANTENNAS CHARACTERISTICS

Model Number GTE AN/18I
Frequency range 4-18 GHz

Gain 6-12 dB

VSWR(maximum) 3.2:1
Isolation <20dB below 5.5 GHz

>20dB above 5.5 GHz

APPENDIX B ANECHOIC CHAMBER

The anechoic chamber is internally lined with absorber material; this material provides the necessary attenuation to the reflection from the walls. Under the absorber material there is an aluminium surface for isolation against external sources of noise such as atmospheric noise and man made noise sources.

Physical dimensions:

Longitudinal length 20 ft.

Lateral length 10 ft.

Height 10 ft.

The absorber material used in the walls of the chamber is EHP-8 Rantec microwave absorber; cut into a precise pyramidal configuration. Figure B.l shows the mechanical specifications as well as the maximum reflections at normal incidence. For the material used in the floor, ceiling, lateral and front walls. The absorber material used for the back wall operates with more absorption at lower frequencies, and the height of the cones is also different. The specifications of the back wall absorber material are shown in Figure B.2.

MECHANICAL SPECIFICATIONS:

| Absorber | Absor | Absorber Height (In.) | (In.) | Pyramid Base | Pyramids |
|------------|---------|-----------------------|---------|--------------|--------------|
| Size (In.) | Overall | Base | Pyromid | Size (In.) | Per Absorber |
| 24×24 | 8-1/2 | 7/1-1 | 7 | 3 × 3 | 64 |

MAXIMUM REFLECTION AT NORMAL INCIDENCE:

| 300 200 120 MH- AMH- | 311111 | | |
|----------------------|--------|-------|--|
| 500 AAH - | 7118 | | |
| 7 | Dillo | 30 db | |
| S | Duba | 40 db | |
|)) | paud | 45 db | |
| × | pana | 50 db | |
| Κυ | pana | 50 db | |

Figure B.1 Specifications of the Absorber Material of the Front Wall.

MECHANICAL SPECIFICATIONS:

| Absorber | Abso | Absorber Height (In.) | nt (In.) | Pyramid Base | Pyramids |
|------------|---------|-----------------------|----------|--------------|--------------|
| Size (In.) | Overall | Bose | Pyramid | Size (In.) | Per Absorber |
| 24 × 24 | 18-1/4 | 2-1/4 | 91 | 9×9 | 91 |

MAXIMUM REFLECTION AT NORMAL INCIDENCE:

| κ | × | U | S | 1 | 200 | 008 | 200 | 120 |
|-------|-------|-------|-------|-------|-------|-----|-----|-----|
| Band | Band | Band | Band | Band | MHz | MHz | MHz | MHz |
| 50 db | 50 db | qp 05 | 45 db | 40 db | 30 db | | | |

Figure B.2 Specifications of the Absorber Material of the Back Wall.

APPENDIX C

SPHERE PROGRAM

```
16
    "SPHERE DRIVES"
  20
30
       COMPUTE BACK SCATTERRED
    I FAR-FIELD FROM A PERFECTLY
  40
      CONDUCTING SPHERE.
 50
       THE INCIDENT FIELD IS A
  SA.
       LINEARLY POLARIZED PLANE W
  70
     AVE WITH ZERO PHASE AT THE CENTER OF THE SPHERE.
       THE THEORETICAL VALUES TO
      BE COMPUTED ARE THE BACK-SC
     ATTERING CROSS-SECTION AND
 90 1 THE PHASE OF THE FAR FIELD
      INTERPOLATED TO THE CENTER
     OF THE SPHERE.
100
       THEORETICAL VALUES ARE
STOPED IN FILES OF 800
110
120
    ! RECORDS, ONE FOR EACH FREQ
UENCY FROM 2.02 GHZ TO 18 GH
130
     Z AT 0.02 GHZ STEPS.
140
150
       "THEOR1. DRIVE1" FOR THE
       1" DIAMETER SPHERE
160
       "THEORY DRIVEL" FOR THE
170
       3 187" DIAMETER SPHERE
180
      "THEOR4 DRIVE1" FOR THE
190
      4.75" DIAMETER SPHERE
200 !
       "THEORE. DRIVE1" FOR THE
210 !
      6" DIAMETER SPHERE
220 !
230
    ! FILE H# STORES THE
    COMPUTED RESULT
250 H#="THEOR6.DRIVE1"
260
270
   -A0=6≭.0254/2 ! SPHERE
    PADIUS IN METERS
280
290 X9=2≭PI ! PARAMETER
300
310 01=2 ! STARTING FREQ IN GHZ
320 02=18 ! FINAL FREQ IN GHZ
330 04=.02 ! FREQ STEP IN GHZ
340 ON ERPOR GOTO 360
350 PURGE H≉
360 OFF ERROR
370 CREATE H$,300,16 ! OPEN A NE
    W FILE WITH 800 RECORDS
380
    ! OF 16 BYTES EACH, EVERY
390
    ! RECORD STRORES ONE MAGNITU
    DE AND ONE PHASE DATA.
```

```
410 ASSIGN# 1 TO H$
420 DIM B8(144), B9(144), D9(144),
    D9(144) ! 144>L0=INT(2*K0*A0
     +3)
430 F0=Q1
440 FOR I=1 TO 800
450 DISP "FREQ LOOP=",I
460 F0=F0+Q4
470 DISP "FREQ (GHZ)=",F0
480 K1=.3/F0 ! WAVELENGTH
490 K0=X9/K1 ! WAVE NUMBER
500 GOSUB 610
510 DISP "E =",E0
520 DISP "P =",P0
530 PRINT# 1,I ; E0,P0
540 NEXT
550 ASSIGN# 1 TO *
560 CLEAR
570 DISP "END OF COMPUTATION"
580 END
590
600
610
620 L0=INT(2*K0*A0+3)
630 IF L0<145 THEN 650
640 DISP "K0*A0 TOO LARGE FOR
    CURRENT ARRAY DIM"
650 Z=K0*A0
660 GOSUB 890
670 E8=0
680 E9=0
690 FOR N=1 TO L0
700 L=L0-N+1
710 M8=D8(L)^2+D9(L)^2
720 M9=B8(L)~2+B9(L)^2
730 A7=(L+.5)/M8/M9
740 A8=A7*(B9(L)*D9(L)-B8(L)*D8(
    L>>
750 A9=A7*(B8(L)*D9(L)+B9(L)*D8(
    LDD
760 E8=A8-E8
770 E9=A9-E9
780 NEXT N
790 E8=-E8
800 E9=-E9
810 E0=E8^2+E9^2
820 P0=ATN2(E9,E8)
830 E0=E0/K0*2*K1 ! CROSS-SECTIO
    Н
840 P0=P0-X9*INT(P0/X9)
950 IF P0<PI THEN 870
860 P0=P0-X9
870 P0=-P0
880 PETURN
890
900 IF Z>L0-1 THEN 1210
910 Z2=Z^2/2
```

```
920 N2=2*Z2+L0+1
930 DI=2*N2+3
940 D2=D1*(2*N2+5)
950 D3=D2*(2*N2+7)
960 D4=T3#(2*N2+9)
970 F1=1-Z2/D1+Z2^2/(2*D2)-Z2^3/
    (6#03)
980 F2=Z*(1/D1-Z2/D2+Z2^2/(2*D3)
    -22^3/(6*04))
990 M=2*Z2
1000 S1=F1
1010 F1=(2*M+1)*F1/Z-F2
1020 F2=S1
1030 IF ABS(F1)(1.E100 THEN 1070
1040 F1=F1*1 E-100
1050 F2=F2*1.E-100
1060 S1=S1*1 E-100
1070 M=M-1
1080 IF M+1>L0 THEN 1000
1090 B8(L0)=F2
1100 BS(L0-1)=F1
1110 NO=L0-2
1120 FOR K=1 TO NO
1130 N=L0-K-1
1140 B8(N)=(2*N+3)*B8(N+1)/Z-B8(
     N+20
1150 NEXT K
1160 A1=(SIN(Z)/Z-COS(Z))/B8(1)
1170 FOR K=1 TO L0
1180 B8(K)=A1*B8(K)
1190 NEXT K
1200 GOTO 1260
1210 B8(1)=SIH(Z)/Z-COS(Z)
1220 B8(2)=(3/Z^2-1)*SIN(Z)-3*CO
     S(Z)/Z
1230 FOR N=3 TO L0
1240 B8(N)=(2*N-1)*B8(N-1)/Z-B8(
     N-2
1250 NEXT N
1260 B9(1)=-SIN(Z)-COS(Z)/Z
1270 B9(2)=(1-3/Z^2)*C0S(Z)-3*SI
     N(2)ノス
1280 FOR N=3 TO L0
1290 B9(N)=(2*N-1)*B9(N-1)/Z-B9(
     N-2)
1300 NEXT N
1310 D8(1)=(1-1/2^2)*SIN(2)+COS(
     Z) / Z
1320 D9(1) = (1/Z^2-1) *COS(Z) + SIN(
     こ)/乙
1330 FOR N=2 TO L0
1340 D8(N)=88(N-1)-N*88(N)/Z
1350 D9(N)=89(N-1)-N*89(N)/Z
1360 NEXT N
1370 RETURN
```

APPENDIX D

CALIB PROGRAM

```
"CALIE DRIVEG"
  20
  30
       CALIBRATION USING A SPHERE
  40
                   GHZ AT F9 GHZ
    1 0VER L9-09
       STEPS BASED ON THEORETICAL
      VALUES COMPUTED USING THE P
     POGRAM "SPHERE DRIVEO"
    ! THE RESULTED SYSTEM TRANS-
     FER FUNCTION IS STORED AS:
       "CALIB3.DRIVE1" (3 187")
 80
 90
    ! A$ IS THE FILE STORING THE
    ! BACKGROUND DATA.
100
    ! C# IS THE FILE STORING THE
     SYSTEM TRANSFER FUNCTION
    ! H$ IS THE FILE STORING
     THEORETICAL DATA OF THE SPHE
    RE
    ! S# DESCRIBES THE SPHERE
130
140
    C#="CALIB3.DRIVE1"
150
160 H$="THEOR3.DRIVE1"
170 S$="3,187 INCH SPHERE"
180 As="BKGRND.DRIVE1"
190 M9=2*PI ! A PARAMETER
200
210 OPTION BASE 1
220 N0=3 ! NUMBER OF READINGS
230 ! TAKEN AND AVERAGED FOR ONE
240 ! FREQ.
250
260 F9=.1 ! FREQ. STEP IN GHZ
270
280 Mi=51 ! Mi=(U9-L9)/F9+2
290
    ! NUMBER OF FREQ. CHECKED.
300
310 L9=10.1 ! LOWER FREQ.IN GHZ
320 U9=15 ! UPPER FREQ.IN GHZ
330
340 DIM A(51,2) ! BACKGROUND
                       DATA
350 DIM B(51,2) ! TARGET DATA
360 DIM G3(51/2) ! THEORY
   ! CREATE C$,M1,16
370
   ! CREATE A$,M1,16
380
   ! STORE CALIBRATION AND BACK
    GROUND DATA IN A FILE OF MI
    RECORDS
   ! EACH RECORD CONTAINS ONE
MAGNITUDE AND ONE PHASE
410 ! DATA AT A FREQUENCY
```

```
420 !
 430 ! READING THE THEORETICAL DA
      TA
 449
 450 ASSIGN# 1 TO H$
 460 K0=(L9-2-F9*2)*50
470 FOR I=1 TO M1
 480 K0=K0+50*F9
 490 READ# 1,K0 ; G3(I,1),G3(I,2)
 500 NEXT I
 510 ASSIGN# 1 TO *
 520 GOSUB 2170 ! HEADER.
530 DISP "DO YOU WANT TO USE THE
       MOST RECENT BACKGROUND DATA
 540 INPUT P#
 550 IF P$="N" THEN 630
 560
      ! READING BACKGROUND DATA.
 570
 580
 590 ASSIGN# 4 TO A$
 600 READ# 4 /
                 \mathbf{a}(\cdot)
    ASSIGN# 4 TO *
 610
 620 GOTO 880
630 CLEAR
640 REMOTE 7 ! REMOTE ALL
                       DEVICES
650 CLEAR 7 ! CLEAR ALL DEVICES
     ! INITIALIZE SIG.GEN TO FIRS
660
     T FREQ.
670 OUTPUT 719 ; "P", L9, "Z1K0L3M0
     N601"
680 CLEAR
690 DISP "REMOVE TARGET FROM CHA
     MBER, PUSH 'CONT' WHEN READY"
700 LOCAL 7
710 BEEP @ BEEP
720 PAUSE
730 REMOTE 7
740 CLEAR
750 DISP "TAKING BACKGROUND DATA
760 PRINT
770 ! PRINT "BACKGROUND DATA"
780 PRINT
790 OUTPUT 719 ; "P", L9, "Z1K0L3M0
     N601"
800 WAIT 200 ! WAIT FOR FREQ.TO
      STABILIZE
810 GOSUB 1320
820
830
       STORING BACKGROUND DATA
849
850 ASSIGN# 3 TO A$
860 PRINT# 3 ; A(,)
870 ASSIGN# 3 TO *
880 CLEAR
```

```
890 LOCAL 7
900 DISP "PUT TARGET INTO CHAMBE
     RIPUSH 'CONT' WHEN READY"
910 DISP "TARGET IS ",S$
920 BEEP @ BEEF
930 PAUSE
940 REMOTE 7
950 CLEAR
960 DISP "COMPUTING TARGET DATA"
970 PRINT
980 ! PRINT "TARGET DATA"
990 PRINT
1000 OUTPUT 719 ; "P", L9, "Z1K0L3M
      9N601"
1010 WAIT 200
1020 GOSUB 1910
1030 PRINT " "
1040 PRÎNT "TRANS. FUNCTION",S$
1050 PRINT " "
1060 !
1070 ! CALCULATE AND STORE TRANS
     FER FUNCTION.
1080
1090 ASSIGN# 2 TO C$
1100 FOR M≈1 TO M1
1110 H1=B(M,1)-A(M,1)
1120 N2=B(M,2)-A(M,2)
1130 X6=G3(M,1)/(N1^2+N2^2)
1140 X7=G3(M. 2)-ATN2(N2,N1)
1150 X7=X7-X9*INT(X7/X9)
1160 IF X7>PI THEN X7=X7-X9
1170 PRINT# 2,M / X6,X7
1180 ! PRINT USING 970 ; M.X6,X7
1190 IMAGE DD.1X, "X6=",SD.DDDE,1
     X, "X7=". SD.DDDE
1200 NEXT M
1210 ASSIGN# 2 TO *
1220 CLEAR
1230 DISP "CALIBRATION COMPLETED
     JDATA STORED IN"/C$
1240 BEEP @ BEEP @ BEEP
1250 LOCAL 7
1260 END
1270
1280
1290
1300
1310
1320
      BACKGROUND DATA COLLECTIO
     N SUBPOUTINE
1330
1340
     ! OUTPUT(L9-F9)TO U9 GHZ AT
      F9
          GHZ STEPS
1350 J=10*(L9-2*F9) ! FREQUENCY
     STARS AT L9-F9 GHZ
```

```
1360 FOR K=1 TO M1 ! NUMBER OF
      FREQUENCY STEPS
 1370 J=J+10#F9
1380 IMAGE 1A,3Z,14A
 1390 OUTPUT 719 USING 1380 ; "P"
      ,J,"00Z1K0L3M0N6O1"
 1400 ! TAKE DATA IN FROM 722%720
1410 GOSUB 1580
1420 !
1430 ! CALCULATE REAL&IMAGINARY
      FROM AMP. &PHASE
1440
1450 R1=A1*COS(P1)
1460 I1=A1*SIN(P1)
1470 A(K,1)=R1
1480 A(K,2)=I1
1490 ! PRINT USING 2040 ; A(K,1)
      ,A(K,2)
1500 NEXT K
1510 OUTPUT 719 / "P",L9, "Z1K0L3M
      0N601"
1520 RETURN
1539
1540
1550
1560
1570
1580 ! SUBROUTINE TO ENTER AMPLI
     TUDE AND PHASE DATA FROM DI
     GITAL VOLTMETERS
1590
1600 ! PREPARE DIGITAL VOLTMETER
      TO SEND AMPLITUDE DATA
1610 ! NO READINGS TAKEN AND AVE
RAGED FOR ONE FREG.
1620
1630
1640 V1=0 ! PARAMETERS FOR THE
1650 F1=0 ! AMERAGING PROCESS.
1660 FOR L=1 TO NO
1670 OUTPUT 720 : "POFIRITIZIFLOM
     Ø "
1680 WAIT 10
1690 ENTER 720 ; V
1700 WAIT 10
1710 OUTPUT 722 ; "F1R7T1M3A0H1"
1720 WAIT 10
1730 ENTER 722 ; F
1740 V1=V1+V
1750 F1=F1+F
1760 WAIT 10
1770 NEXT L
1780 V=V1/N0
1790 F=F1/N0
1800 A1=10^F ! TRANSFER TO MAG
      FROM VOLTS
```

```
1810 P1=100*V ! TRANSFER TO DEG.
      FROM VOLTS.
 1820 P1=DTR(P1)
       PRINT USING 1810 ; KJA1,P
 1830
 1840 IMAGE 00/2%,"A=",MD.000E,2%
      "P=" SD DDDE
1850 RETURN
 1860
1870
1880
1890
1900
1910
       TARGET DATA COLLECTION
       SUBROUTINE
1929
        OUTPUT(L9-F9)TO U9 GHZ AT
1930
      F9 GHZ STEPS
      J=10*(L9-2*F9) ! FREQUENCY
1940
      STARS AT L9-F9 GHZ
1950 FOR K=1 TO M1 ! NUMBER OF
     FREQUENCY STEPS
1960
     J=J+10*F9
1970 IMAGE 1A,3Z,14A
1980 OUTPUT 719 USING 1970 ;
      ,J,"00Z1K0L3M0N601"
1990
     ! TAKE DATA IN FROM 722%720
     1820 GOSUB 1410
2000 GOSUB 1580
2010 ! CALCULATE REAL&IMAG.FROM
     AMP&PHASE
2020 R1=A1*COS(P1)
2030 I1=A1*SIN(P1)
2040 B(K,1)=R1
2050 B(K,2)=11
2060 ! PRINT USING 2040 ; B(K,1)
      /B(K/2)
2070 IMAGE 4%,"R=",SD.DDDE,2%,"I
     =".SD.DDDE
2080 NEXT K
2090 OUTPUT 719 ; "P", L9, "Z1K0L3M
     ØN601"
2100 RETURN
2110
2120
2130
2140
2150
     ! HEADER SUBROUTINE
2160
2170 PRINT
2180 PRINT
2190 CLEAR
    DISP "CALIBRATION STANDARD"
2200
      9.≴
2210 PRINT "CALIBRATION STANDARD
     " » S$
```

APPENDIX E

TARGET PROGRAM

```
"TARGET DRIVEG"
  10
  20
  30
       TARGET BACK-SCATTERING
       USING C$ DATA & STORE
  40
       RESULTS IN G#
     ! FREQUENCIES: L9-U9 GHZ AT
  50
      F9
          GHZ STEPS
  60
       FILE C# STORES THE SYSTEM
  70
      TRANSFER FUNCTION
  80 ! FILE G$ STORES TARGET DATA
      OBTAINED FROM THIS PROGRAM
 90 ! FILE H# STORES THEORETICAL VALUES FOR PLOTTING OVERLAY
100 ! FILE A$ STORES BACKGROUND
      DATA
110 C$="CALIB3.DRIVE1"
120 G$="SCR3.DRIVE1"
130 H#="THEOR3.DRIVE1"
140 A$="BKGRND.DRIVE1"
159
169
       CREATE G$,52,24
       STORE TARGET DATA IN FILE
170
       OF M1+1 RECORDS FIRST ONE
180
199
       FOR THE AVERAGE PROCEDURE
200
     ! AND THE REST CONTAINS THE
210
       FREQUENCY MAGNETUDE AND
220
     ! PHASE SHIFT.
230
240
      CREATE A$,51,16
    ! STORE CALIB. AND BACKGROUN
DATA IN A FILE OF M1 RECORDS
250
     ! EACH RECORD CONTAINS ONE M
    AG. AND PHASE AT A FREQ.
279
280 OPTION BASE 1
290 NO=2 ! NUMBER OF READINGS
300 ! TAKEN AND AVERAGED FOR ONE
        FREQUENCY.
310
320 DIM A(51,2) ! BACKGROUND DAT
330 DIM B(51,2) ! TARGET DATA
340 DIM G4(51,2) ! CALIBRATION
350 DIM N(51,3) ! RESULTANT
360 DIM M9(51/3)
370
380 N1=51 !
390 ! N1=(U9-L9)/F9+2 NUMBER OF
400 | FREQ. CHECKED.
410 F9=.1 ! FREQ.STEPS IN GHZ.
```

```
420 U9=15 ! UPFER FREQ. IN GHZ
430 L9=10.1 ! LOWER FREQ. IN GHZ
440 DIM T(800,2) ! STORES THEORE
      TICAL DATA
 450 X9=2*PI
 460
 470
      ! READING TRANSFER FUNCTION
 480
 490 ASSIGN# 1 TO C$
 500 READ# 1 / G4(/)
510 ASSIGN# 1 TO *
 520 ! MAT PRINT USING 330 ; G4
 530 IMAGE 2%,30.40
 540 !
 550 REMOTE 7 ! REMOTE ALL DEVICES
 560 CLEAR 7 ! CLEAR ALL DEVICES
570 OUTPUT 719 ; "P1Z1K0L3M0N601"
       ! INITIAL SETUP OF 719
 580 CLEAR
 590 DISP "DO YOU WANT TO USE THE
      MOST RECENT BACKGROUND DATA
              YZN"
 600 INPUT P#
610
620 IF P$="N" THEN 700
630
640 ! READING BACKGROUND DATA
650 !
660 ASSIGN# 4 TO A$
670 READ# 4 / A(/)
680 ASSIGN# 4 TO *
690 GOTO 860
700 DISP "REMOVE TARGET FROM
CHAMBER, PUSH 'CONT' WHEN REA
     BY"
710 LOCAL 7
720 BEEP @ BEEP
730 PAUSE
740 DISP "TAKING BACKGROUND DATA
750 REMOTE 7
760 OUTPUT 719 ; "P", L9, "Z1K0L3M0
     N601" ! INITIAL SETUP OF 719
770 WAIT 100
780 GOSUB 2520
790 !
    ! STORING BACKGROUND DATA.
800
819
820 ASSIGN# 5 TO A$
830 PRINT# 5 ; A(,)
840 ASSIGN# 5 TO *
850 CLEAR
860 DISP "PUT TARGET INTO CHAMBE
     RIPUSH 'CONT' WHEN READY"
870 LOCAL 7
880 BEEP @ BEEP
890 PAUSE
```

•

```
900 REMOTE 7
910 OUTPUT 719 : "P", L9, "Z1K0L3M0
     N601" ! INITIAL SETUP OF 719
 920 WAIT 500
 930 GOSUB 3370
 940 CLEAR
 950 DISP "COMPUTING TARGET DATA"
 960 GOSUB 3090
 970
 980 !
       COMPUTING TARGET DATA
 990 !
       WITHOUT BACKGROUND AND THE
 1000 ! FREQ. FOR EACH RECORD.
1010 F0=L9-2*F9
1020 FOR M=1 TO N1
1030 F0=F0+F9
1040 H(M,1)=F0
1050 X7=B(M,1)-A(M,1)
1060 X8=B(M,2)-A(M,2)
1070 X6=(X7^2+X8^2)*G4(M,1)
1080 N(M,2)=X6
1090 X8=ATN2(X8,X7)+G4(M,2)
1100 X8=X8-X9*INT(X8/X9)
1110 IF X8>PI THEN X8=X8-X9
1120 N(M,3)=X8
1130 NEXT M
1140 DISP "PRINT DATA? YZN"
1150 BEEP @ BEEP
1160 INPUT P#
1170 IF P$="N" THEN 1210
1180 PRINT "
                  FREQ
                             CRSEC
          PHASE"
1190 MAT PRINT USING 1200 ; N
1200 IMAGE 2X/3D.4D
1210 CLEAR
1220 LOCAL 7
1230 DISP "PLOT MAGNITUDE FOR
           THIS MEASURMENT? YZN"
1240 INPUT P≴
1250 IF P$="H" THEN 1290
1260 DISP "SELECT PEN. PUSH
          'CONT' WHEN READY"
1270 PAUSE
1280 GOSUB 3530
1290 CLEAR
1300 DISP "PLOT PHASE FOR THIS
          MEASURMENT ? YZN"
1310 BEEP @ BEEP
1320 INPUT P#
1330 IF P$="N" THEN 1370
1340 DISP "SELECT PEN. PUSH
           'CONT' WHEN READY"
1350 PAUSE
1360 GOSUB 4570
1370 CLEAR
```

```
1380 DISP "DO YOU WANT TO MAKE
            AVERAGE WITH PREVIUSE
       DATA?"
 1390 DISP "2Y/N"
 1400 BEEP @ BEEP
 1410 INPUT P$
1420 IF P$="Y" THEN 1710
1430 DISP "DO YOU WANT TO STORE
            DATA ? Y/N "
 1440 INPUT P$
 1450 IF P$="H" THEN 2240
1460 MO=1
           "DO YOU WANT TO STORE
1470 DISP
            DATA IN FILE"
1480 DISP G$
 1490 DISP "? Y/N"
1500 INPUT P$
1510 IF P$="Y" THEN 1580
1520 DISP "ENTER NAME OF THE DAT
      A FILE TO BE USED FOR STORE
      GE"
1530 INPUT G#
1540 DISP "IS THIS AN OLD FILE
        TO BE UPDATED ? YZN "
1550 INPUT P$
1560 IF P$="Y" THEN 1580
1570 CREATE G$,53,24
1580 DISP "ENTER LENGTH OF TARGE
1590 BEEP @ BEEP
1600 INPUT MI
1610 DISP "ENTER DIAMETER OF TAR
      GET"
1620 BEEP @ BEEP
1630 INPUT M2
1640
1650 ! STORE MEASURED DATA.
1660
1670 ASSIGN# 2 TO G$
1680 PRINT# 2 ; M0/M1/M2/N())
1690 ASSIGN# 2 TO *
1700 GOTO 2240
1710 DISP "DOES THE DATA STORED
           IN FILE"
1720 DISP G#
1730 DISP "? Y/N"
1740 INPUT P$
1750 IF P$="Y" THEN 1830
1760 DISP "ENTER NAME OF DATA
          FILE TO BE USED FOR THE
           AVERAGE"
1770 !
1780 INPUT G$
1790 !
1800 ! READ OLD DATA
1810 ! AND MAKES WIGHTED AVERAGE
1820 ! WITH NEW DATA.
```

```
1830 ASSIGN# 6 TO G$
 1840 READ# 6 : MO.M1.M2.M9(,)
 1850 ASSIGN# 6 TO *
 1860 FOR K=1 TO N1
 1870 M9(K,2)=M9(K,2)*M0+N(K,2)
 1880 M9(K,2)=M9(K,2)/(M0+1)
 1890 M9(K,3)=M9(K,3)*M0+N(K,3)
1900 M9(K,3)=M9(K,3)/(M0+1)
 1910 N(K,2)=M9(K,2)
 1920 N(K,3)=M9(K,3)
1930 NEXT K
 1940 M0=M0+1
 1950
 1960 !
         STORE NEW AVERAGE.
 1970 ASSIGN# 7 TO G$
 1980 PRINT# 7; M0,M1,M2,N(,)
1990 ASSIGN# 7 TO *
2000 PRINT "DATA IS AVERAGE OF",
       MO . "MEASURMENTS"
 2010 DISP "PRINT DATA? Y/N"
 2020 BEEP @ BEEP
 2030 INPUT P$
 2040 IF P$="N" THEN 2080
2050 PRINT "
                     FREQ
                                  CRSEC
            PHASE"
2060 MAT PRINT USING 1200; N
2070 IMAGE 2%; 3D.4D
2080 DISP "PLOT MAGNITUDE? Y
2090 BEEP @ BEEP
2100 INPUT P$
2110 IF P$="H" THEN 2150
2120 DISP "SELECT PEN FOR MAGNIT
      UDE PLOT. PUSH 'CONT' WHEN
      READY.
2130 PAUSE
2140 GOSUB 3530
2150 CLEAR
2160 DISP "PLOT PHASE? Y/N"
2170 BEEP @ BEEP
2180 INPUT P#
2190 IF P$="N" THEN 2230
2200 DISP "SELECT PEN AND CHANGE
PAPER FOR PHASE PLOT. PUS
      H 'CONT' WHEN READY."
2210 PAUSE
2220 GOSUB 4570
2230 CLEAR
2240 DISP "DO YOU WANT TO"
2250 DISP "OBTAIN DATA"
2260 DISP "FOR A NEW TARGET?"
2280 DISP "ENTER Y/N"
2290 INPUT P*
2300 IF P$="N" THEN 2430
2310 DISP "DO YOU WANT TO USE
THE SAME FILE"
2320 DISP G#
```

```
2330 DISP "TO STORE NEW DATA? Y/
      H"
2340 INPUT P$
2350 IF P$="Y" THEN 2420
2360 DISP "ENTER NEW FILE NAME T
      O STORE TARGET DATA"
2370 INPUT G$
2380 DISP "IS THIS AN OLD FILE
     TO BE UPDATED? Y/N"
2390
 2400 IF P$="Y" THEN 2420
2410 CREATE G$,52,24
 2420 GOTO 550
2430 CLEAR
2440 DISP "END OF PROGRAM"
2450 BEEP @ BEEP @ BEEP
2460 END
 2470
2480
2490
2500
2510
     ! BACKGROUND DATA COLLECTIO
      N SUBROUTINE
2530
     ! OUTPUT(L9-F9)TO U9 GHZ
      J=10*(L9-2*F9) ! FREQUENCY
      STARTS AT L9-F9 GHZ TO BE
      INCREASED AT F9 GHZ STEPS
2550 FOR K=1 TO N1 ! NUMBER OF F
      REQUENCY STEPS
2560 J=J+10*F9
2570 IMAGE 1A,3Z,14A
2580 OUTPUT 719 USING 2570 ; "P"
      , J. "00Z1K0L3M0N601"
     ! 50 MSEC WAIT FOR FREQUENC
     Ý TÖ STABILIZE
2600 WAIT 50
2610 ! TAKE DATA IN FROM 722 AND
       720
2620 GOSUB 2780
2630 ! REAL AND IMAGINARY PARTS
2640 ! FROM AMP. AND PHASE.
2650 R1=A1*COS(P1)
2660 I1=A1*SIN(P1)
2670 A(K,1)=R1
2680 A(K,2)=I1
2690 ! PRINT "I1=",A(K,2)
2700 ! PRINT "R1=",A(K,1)
2710 NEXT K
2720 OUTPUT 719 ; "P" | L9, "Z1K0L3M
0N601" ! INITIAL SETUP OF 7
     19
2730 RETURN
2740
2750
```

```
2760
2770
 2780
        SUBROUTINE TO ENTER AMPLI
      TUDE AND PHASE DATA FROM DI
      GITAL VOLTMETER
 2790
 2800
      ! PREPARE DIGITAL VOLTMETER
       TO SEND AMPLITUDE DATA
 2810 ! NO READINGS TAKEN AND AVE RAGED FOR ONE FREQUENCY.
 2820 V1=0 ! PARAMETERS FOR AVERA
                  GING PROCESS.
2830 W1=0
2840 FOR L=1 TO NO
2850 OUTPUT 720 ; "F1R1T1Z1FL0M0"
 2860 WAIT 10
2870 ENTER 720 ; VO
2880 WAIT 10
2890 OUTPUT 722 ; "F1R7T1M3A0H1"
2900 WAIT 10
2910 ENTER 722 ; WO
2920 V1=V1+V0
2930 W1=W1+W0
2940 WAIT 10
2950 NEXT L
2960 V0=V1/N0
2970 W0=W1/N0
2980 ! TRANSFERS FROM VOLTS TO
           AMPL.
2990 A1=10^W0
3000 ! TRANSFERS TO DEG. FROM VO
           LTS
3010 P1=100*V0
3020 ! PRINT "A1=",A1
3030 P1=DTR(P1)
3040 ! PRINT "P1=",P1
3050 RETURN
3060
3070
3080
3090
        DATA COLLECTION SUBROUTIN
     ! OUTPUT(L9-F9)TO U9 GHZ AT
3100
     F9 GHZ STEPS
J≈10*(L9-2*F9) ! INI
EQUENCY AT L9-F9 GHZ
3110
                      ! INITIAL FR
3120 FOR K=1 TO N1 ! FREQUENCY S
     TEPS
     U=U+10*F9 ! F9 GHZ INCREME
3130
     NTŚ
3140 IMAGE 1A,32,14A
3150 OUTPUT 719 USING 3140 ;
      ,J,"00Z1K0L3M0N6O1"
     ! 50 MSEC WAIT FOR FREQUENC
     Y TO STABILIZE
3170 WAIT 50
```

```
3180 ! TAKE DATA IN FROM 722%720
3190 GOSUB 2780
3200 !
3210 ! REAL&IMAG FROM AMP.&PHASE
 3220 P1=A1*COS(P1)
3230 I1=A1#SIN(P1)
3240 B(K,1)=R1
3250 B(K,2)=I1
3260 ! PRINT "R1=",B(K,1)
3270 ! PRINT "I1=",B(K,2)
3280 NEXT K
3290 OUTPUT 719 ; "P",L9, "Z1K0L3M
      @N601" ! INITIAL SETUP OF 7
      19
3300 RETURN
3310
3320
3330
3340 ! HEADER SUBROUTINE
3350 !
3360 D$="MONTH/DATE/YEAR"
3370 PRINT " "
3380 PRINT
3390 CLEAR
3400 DISP "ENTER TODAY'S DATE -
     MONTH, DATE, YEAR"
3410 INPUT D#
3420 DISP "ENTER TGT DESCRIPTION
3430 INPUT T#
3440 PRINT D#
3450 PRINT "TARGET IS ",T$
3460 PRINT "*************
3470 PRINT "*************
3480 PRINT
3490 CLEAR
3500 RETURN
3510
3520
3530 ! MAGNITUDE PLOTTING
3540 !
       SUBROUTINE
3550
3560 PLOTTER IS 705
3570 LOCATE 32,122,20,85
3580 FRAME
3590
       SEARCH FOR MAX. & MIN.
3600
       S0=N(2,2)
3610
       $1=90
3620
     ! FOR M=3 TO N1
3630
     ! IF SO>N(M,2) THEN SO=N(M,
     2)
3640
      IF S1<N(M,2) THEN S1=N(M,
    2)
3650
       HEXT M
3660 DISP "ENTER LOWER VALUE FOR
          MAGNITUDE PLOTTING"
```

```
3670 BEEP @ BEEP
3680 INPUT S0
 3690 DISP "ENTER UPPER VALUE FOR
            MAGNITUDE PLOTTING"
 3700 INPUT S1
 3710 L1=INT(L9)
 3720 U1=CEIL(U9)
 3730
 3740 ! CALCULATE SCALE STEPS
 3750 ! FOR MAGNETUDE.
3760 S3=LGT(S1)
 3770 S4=INT(S3)-1
 3780 S5=S3-S4
3790 S5=INT(10^S5)+1
3800 S4=10^S4
3810 L0=INT(S0/S4)
3820 IF S5-L0<=14 THEN 3900
3830 IF S5-L0>=50 THEN 3870
3840 S5=.5*S5
3850 $4=$4*2
3860 GOTO 3810
3870 S5= 2*S5
3880 S4=5*S4
3890 GOTO 3810
3900 L0=S4*L0
3910 U0=S5*S4
3920 D0=U0-L0
3930 SCALE L1,U1,L0,U0
3940 FXD 0,4
3950 LAXES -1,84,L1,L0
3960 MOVE L1,0
3970 FOR K=2 TO N1
3980 W9=N(K,1)
3990 R9=N(K,2)
4000 GOSUB 4470
4010 NEXT K
4020 M5=(U1+L1)/2
4030 MOVE M5.L0-.09*D0
4040 LORG 5 @ CSIZE 3,1,0
4050 ! LABEL "FREQUENCY (GHZ)"
4060 MOVE L1-1, 5*(L0+U0)
4070 LDIR PI/2
4080 ! LABEL "CROSS SECTION (SQ.
     METERS!"
4090 MOVE M5, U0+.09*D0
4100 LDIR 0
4110 CSIZE 3,1,0
4120 LABEL T$
4130 MOVE M5, U0+.03*D0
4140 LABEL D#
4150 PEHUP
4160 DISP "OVERLAY THEORETICAL C
     URVE? YZN"
4170
4180 INPUT P$
4190 IF P$="N" THEN 4450
```

```
4200 DISP "IS THE THEORETICAL
DATA STORED IN THE FILE", H$
4210 DISP "? Y/N"
4220 INPUT P$
4230 IF P$="Y" THEN 4260
4240 DISP "ENTER NAME OF THE
      DATA FILE TO BE PLOTTED."
4250 INPUT H$
4260 BEEP @ BEEP
4270 DISP "CHANGE PEN IF DESIRED . PUSH 'CONT' WHEN READY."
4280 PAUSE
4290 ASSIGN# 3 TO H$
4300 J1=(L9-2)*50
4310 J2=(U9-2)*50
4320 FOR J=J1 TO J2
4330 READ# 3,J ; T(J,1),T(J,2)
4340 NEXT J
4350 ASSIGN# 3 TO *
4360 F0=L9
4370 R9=T(J1,1)
4380 MOVE F0,R9
4390 FOR I=J1+1 TO J2
4400 F0=F0+.02
4410 R9=T(I,1)
4420 DRAW F0,R9
4430 NEXT I
4440 PENUP
4450 RETURN
4460 !
4470 ! PLOT CROSS
4480 MOVE M9.R9
4490 CSIZE 2,.5,0
4500 LABEL "+"
4510 ! IMOVE .00025,.00025
4520 ! IDRAW -.0005,0
4530 RETURN
4540
4550
4560 !
4570 ! PHASE PLOTTING SUBROUTINE
4580 !
4590 PLOTTER IS 705
4600 LOCATE 32,122,20,85
4610 FRAME
4620 S4=.25
4630 U0=1
4640 L0=-1
4650 D0=U0-L0
4660 SCALE L1,U1,L0,U0
4670 FXD 0,3
4680 LAXES -1,54,L1,L0
4690 MOVE L1,0
4700 FOR K=2 TO N1
4710 M9=N(K,1)
4720 R9=N(K,3)/PI
4730 GOSUB 5070
```

```
4740 NEXT K
 4750 MOVE M5,L0-.09*D0
4760 LORG 5 @ CSIZE 3,1,0
 4770 ! LABEL "FREQUENCY (GHZ)"
 4780 MOVE L1-1, 5*(L0+U0)
 4790 LDIR PI/2
 4800 LABEL "PHASE (PI)*
 4810 MOVE M5,U0+.09*D0
 4820 LDIR 0
 4830 CSIZE 3,1,0
4840 LABEL T$
 4850 MOVE M5, U0+.03*D0
 4860 LABEL D$
4870 PENUP
4880 DISP "OVERLAY THEORETICAL C
       URVE? YZN"
4890 INPUT P#
4900 IF P$="N" THEN 5030
4910 BEEP @ BEEP
4920 DISP "CHANGE PEN IF DESIRED . PUSH 'CONT' WHEN READY."
4930 PAUSE
4940 F0=L9
4950 R9=T(J1,2)/PI
4960 MOVE F0,R9
4970 FOR I=J1+1 TO J2
4980 F0=F0+.02
4990 R9=T(I,2)/PI
5000 DRAW F0,89
5010 NEXT I
5020 PENUP
5030 RETURN
5040
5050
5060
5070 ! PLOT DOT
5080 MOVE M9, R9
5090 CSIZE 2, 5,0
5100 LABEL "*"
5110 ! IMOVE .00025,.00025
5120 ! IDRAW -.0005,0
5130 RETURN
```

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